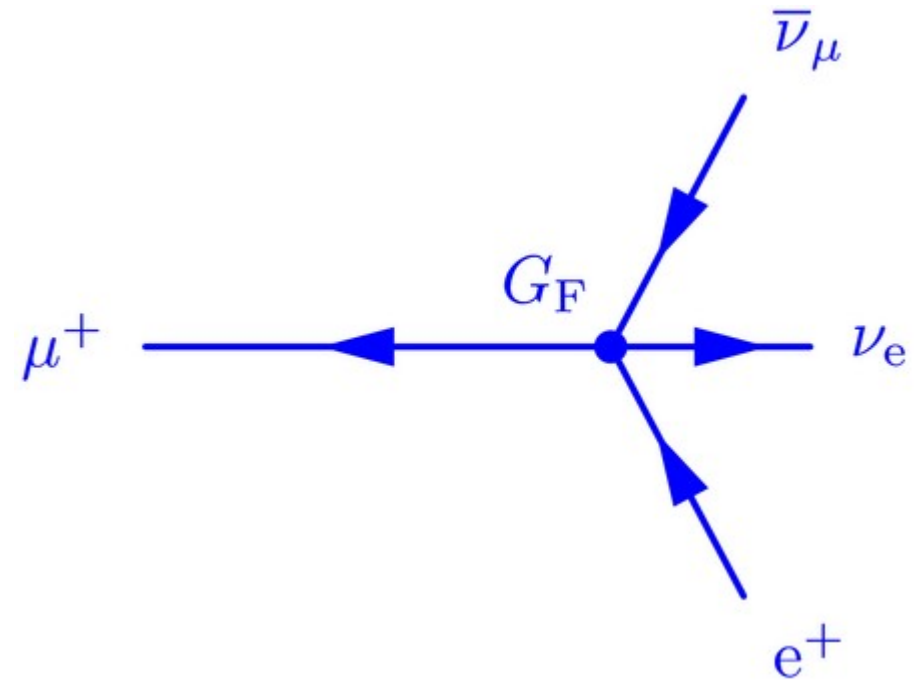


The MuLan Experiment

Measuring the muon lifetime to 1ppm

Kevin Lynch
MuLan Collaboration
York College, CUNY



Berkeley, Boston, Illinois,
James Madison, Kentucky, KVI, PSI

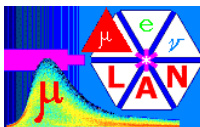
The MuLan Experiment

Measuring the muon lifetime to 1ppm

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York College, CUNY

Outline:
Motivate the measurement
Describe the experiment
Results

Berkeley, Boston, Illinois,
James Madison, Kentucky, KVI, PSI

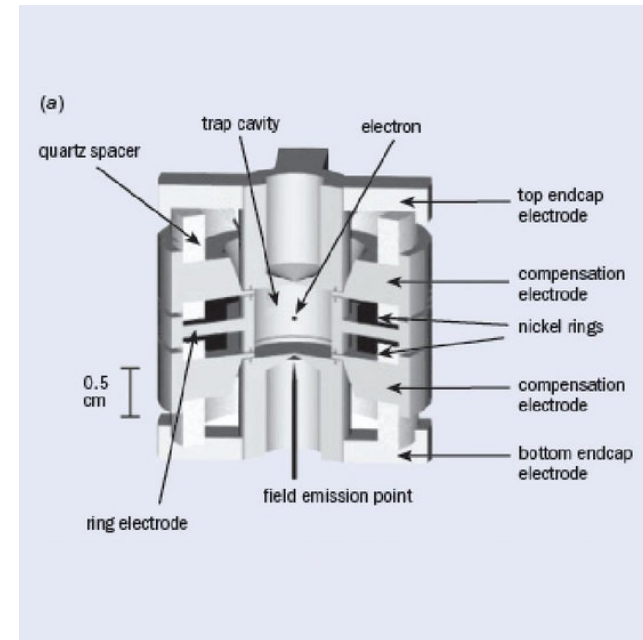


Precision electroweak predictions rest on three input parameters

Fine Structure Constant

$$\frac{\delta\alpha_{\text{em}}}{\alpha_{\text{em}}} \approx 0.32 \text{ ppb}$$

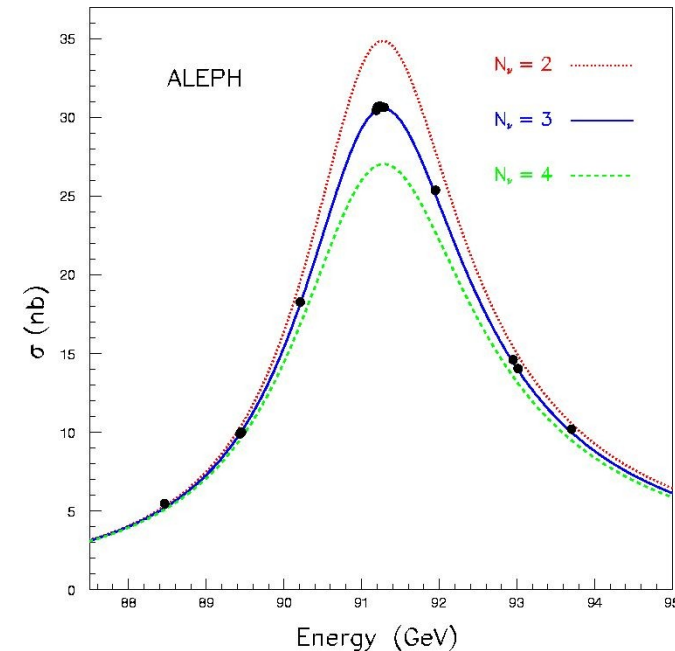
Gabrielse *et al*
2008



Mass of the neutral weak boson

$$\frac{\delta M_{Z^0}}{M_{Z^0}} \approx 23 \text{ ppm}$$

LEP EWWG
2005



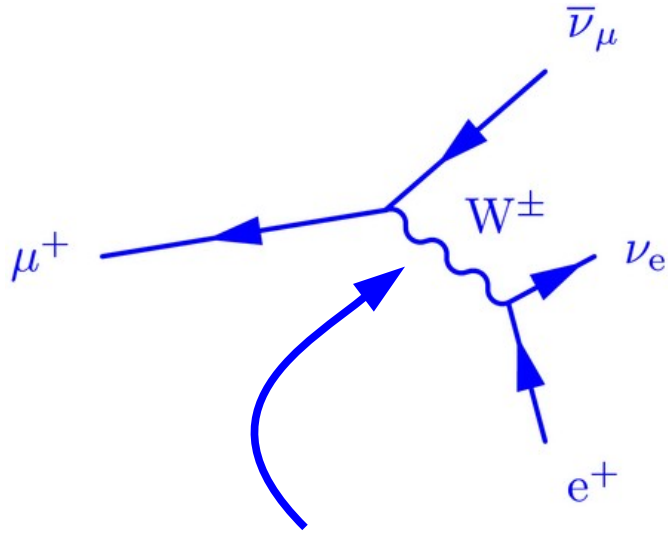
Fermi Constant

$$\frac{\delta G_F}{G_F} \approx 9 \text{ ppm}$$

Giovanetti *et al*
1984

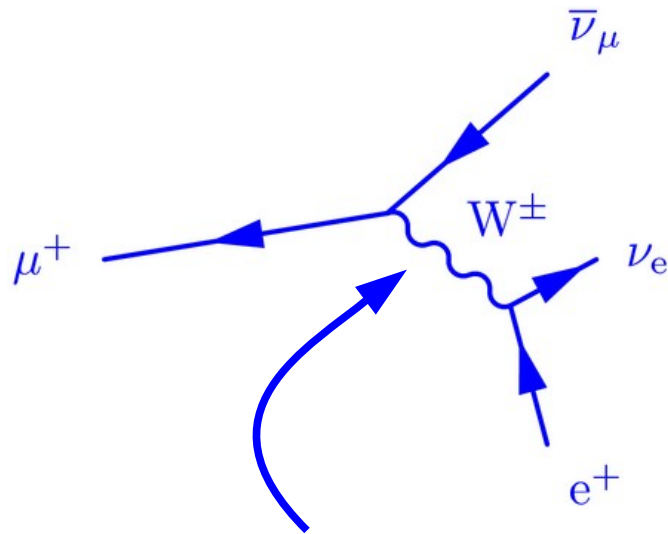
Muon decay gives us unique access to the electroweak scale

Muon decay gives us unique access to the electroweak scale



The muon *only* decays via the weak interaction, which gives it a very long lifetime.

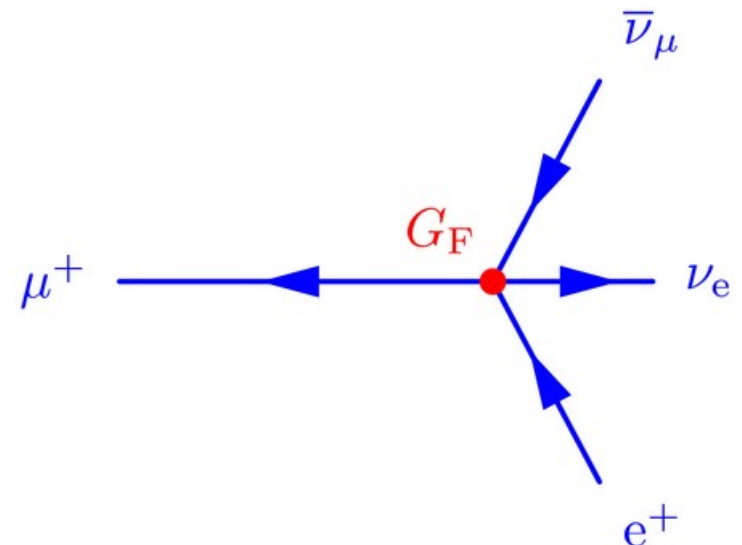
Muon decay gives us unique access to the electroweak scale



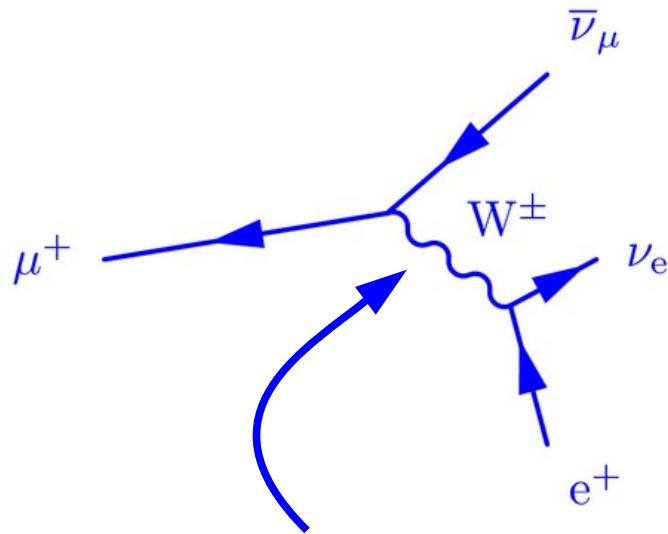
The muon *only* decays via the weak interaction, which gives it a very long lifetime.

The V-A theory factorizes into a pure **weak** contribution

$$\frac{1}{\tau_{\mu^+}} = \frac{G_F^2 m_\mu^5}{192\pi^3}$$



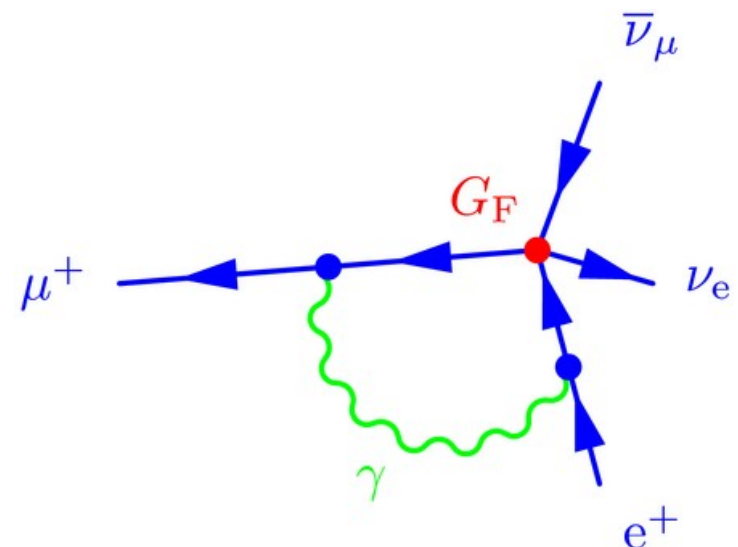
Muon decay gives us unique access to the electroweak scale



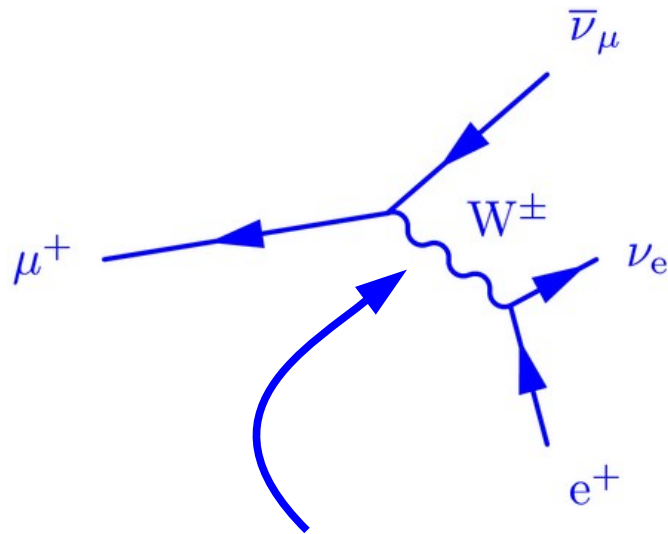
The muon *only* decays via the weak interaction, which gives it a very long lifetime.

The V-A theory factorizes into a pure **weak** contribution, and **non-weak** corrections, essentially uncontaminated by hadronic uncertainties.

$$\frac{1}{\tau_{\mu^+}} = \frac{G_F^2 m_\mu^5}{192\pi^3} (1 + q)$$



Muon decay gives us unique access to the electroweak scale

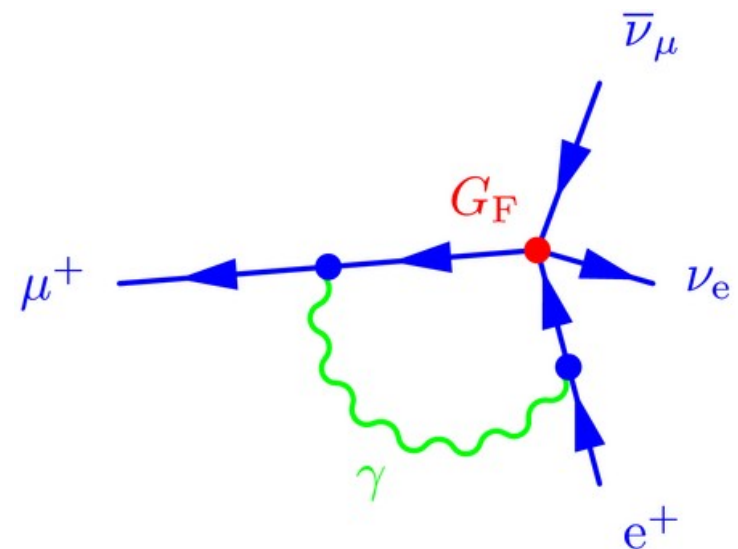


The muon *only* decays via the weak interaction, which gives it a very long lifetime.

All relevant weak interaction physics is confined to one easily measured parameter with a clean theoretical interpretation.

The V-A theory factorizes into a pure **weak** contribution, and **non-weak** corrections, essentially uncontaminated by hadronic uncertainties.

$$\frac{1}{\tau_{\mu^+}} = \frac{G_F^2 m_\mu^5}{192\pi^3} (1 + q)$$

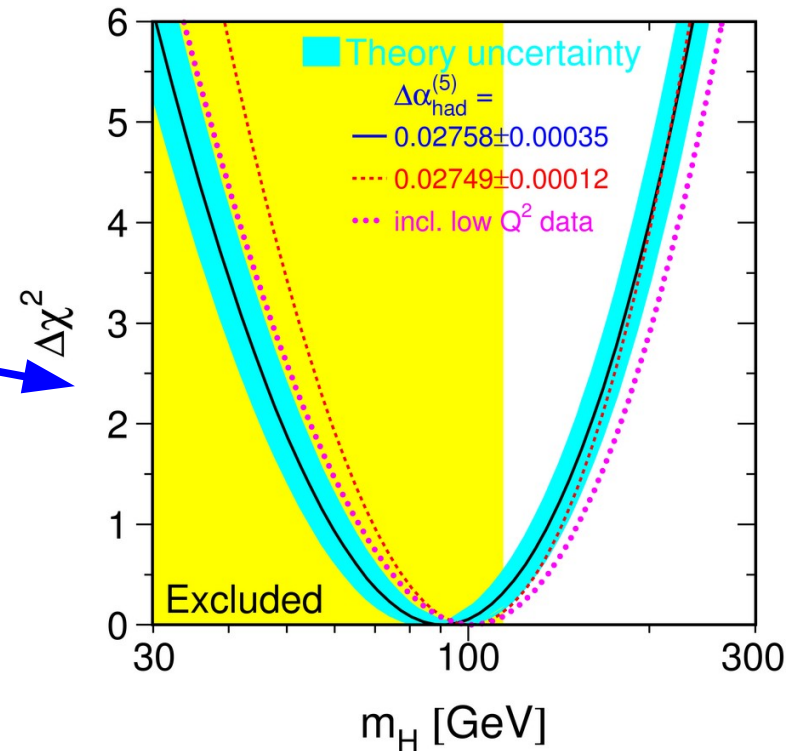


The Fermi constant is an implicit input to all precision electroweak studies

$$\frac{G_F}{\sqrt{2}} = \frac{g^2}{8M_W^2} (1 + \Delta r(m_t, m_H, \dots))$$

Contains all weak interaction loop corrections.

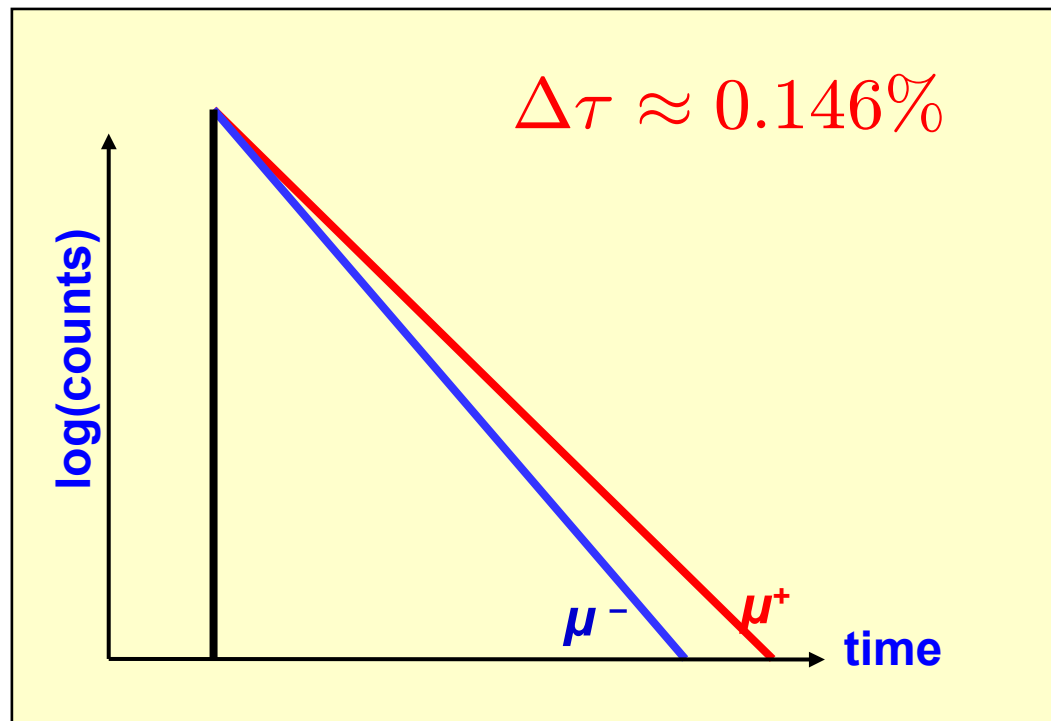
Example: the “blue band” Higgs limit plot.



Precision lifetime difference measurements yield information on nucleon weak structure

For example, the singlet capture rate on the proton gives direct access to the pseudoscalar nucleon coupling

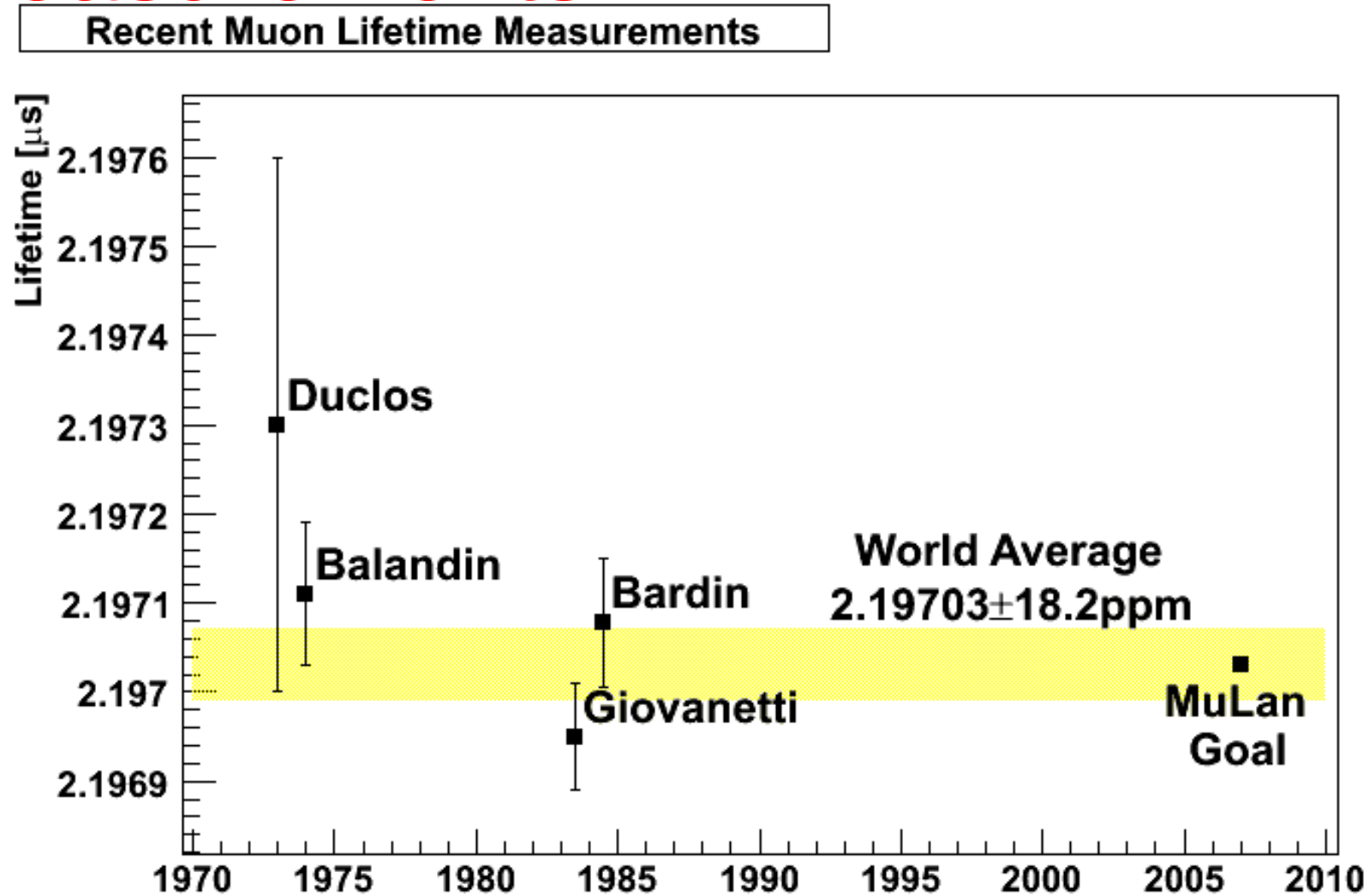
$$\Lambda_S = \Lambda_{\mu^-} - \Lambda_{\mu^+} = \frac{1}{\tau_{\mu^-}} - \frac{1}{\tau_{\mu^+}}$$



$$\frac{\delta\Lambda_S}{\Lambda_S} = 0.18 \frac{\delta g_P}{g_P}$$

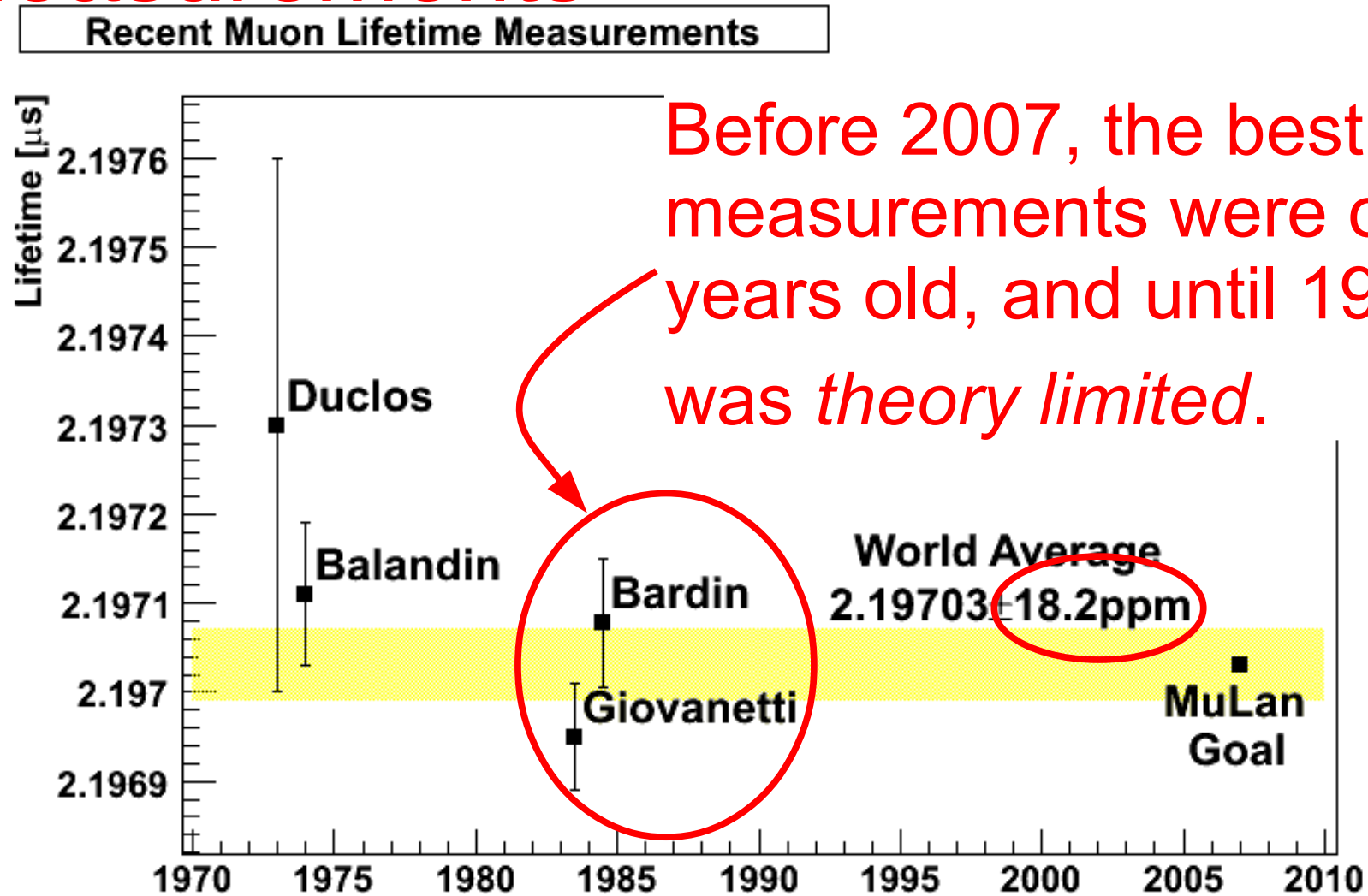
The more accurately we measure the muon lifetimes, the more precisely we can extract derived quantities

A brief history of muons lifetime measurements



G. Bardin et al., Phys. Lett. B 137, 135 (1984)
K. Giovanetti et al., Phys. Rev. D 29, 343 (1984)

A brief history of muons lifetime measurements



G. Bardin et al., Phys. Lett. B 137, 135 (1984)

K. Giovanetti et al., Phys. Rev. D 29, 343 (1984)

Theory limitations were lifted in 1999

$$\frac{\delta G_F}{G_F} = \frac{1}{2} \sqrt{\left(\frac{\delta\tau_\mu}{\tau_\mu}\right)^2 + \left(5\frac{\delta m_\mu}{m_\mu}\right)^2 + \left(\frac{\delta\text{theory}}{\text{theory}}\right)^2}$$

Mid 90s: 17 ppm 18 ppm 90 ppb 30 ppm

Theory limitations were lifted in 1999

$$\frac{\delta G_F}{G_F} = \frac{1}{2} \sqrt{\left(\frac{\delta\tau_\mu}{\tau_\mu}\right)^2 + \left(\frac{\delta\text{theory}}{\text{theory}}\right)^2}$$

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1999:

9 ppm

18 ppm

< 0.3 ppm

van Ritbergen and Stuart:
2-loop QED corrections
(massless electrons)

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$$\frac{\delta G_F}{G_F} = \frac{1}{2} \sqrt{\left(\frac{\delta \tau_\mu}{\tau_\mu}\right)^2 + \left(\frac{\delta \text{theory}}{\text{theory}}\right)^2}$$

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Lifetime error limits
the Fermi constant
extraction

van Ritbergen and Stuart:
2-loop QED corrections
(massless electrons)

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$$\frac{\delta G_F}{G_F} = \frac{1}{2} \sqrt{\left(\frac{\delta \tau_\mu}{\tau_\mu}\right)^2 + \left(\frac{\delta \text{theory}}{\text{theory}}\right)^2}$$

Today:

0.5 ppm

1 ppm

< 0.3 ppm

Lifetime error limits
the Fermi constant
extraction

How do you measure the muon lifetime?

How do you measure the muon lifetime?



$$\tau_{\mu} = 2.197 \mu\text{s}$$

How do you measure the muon lifetime?

One-at-a-time



$$\tau_{\mu} = 2.197 \mu\text{s}$$

How do you measure the muon lifetime?

One-at-a-time



$$\tau_{\mu} = 2.197 \mu\text{s}$$

How do you measure the muon lifetime?

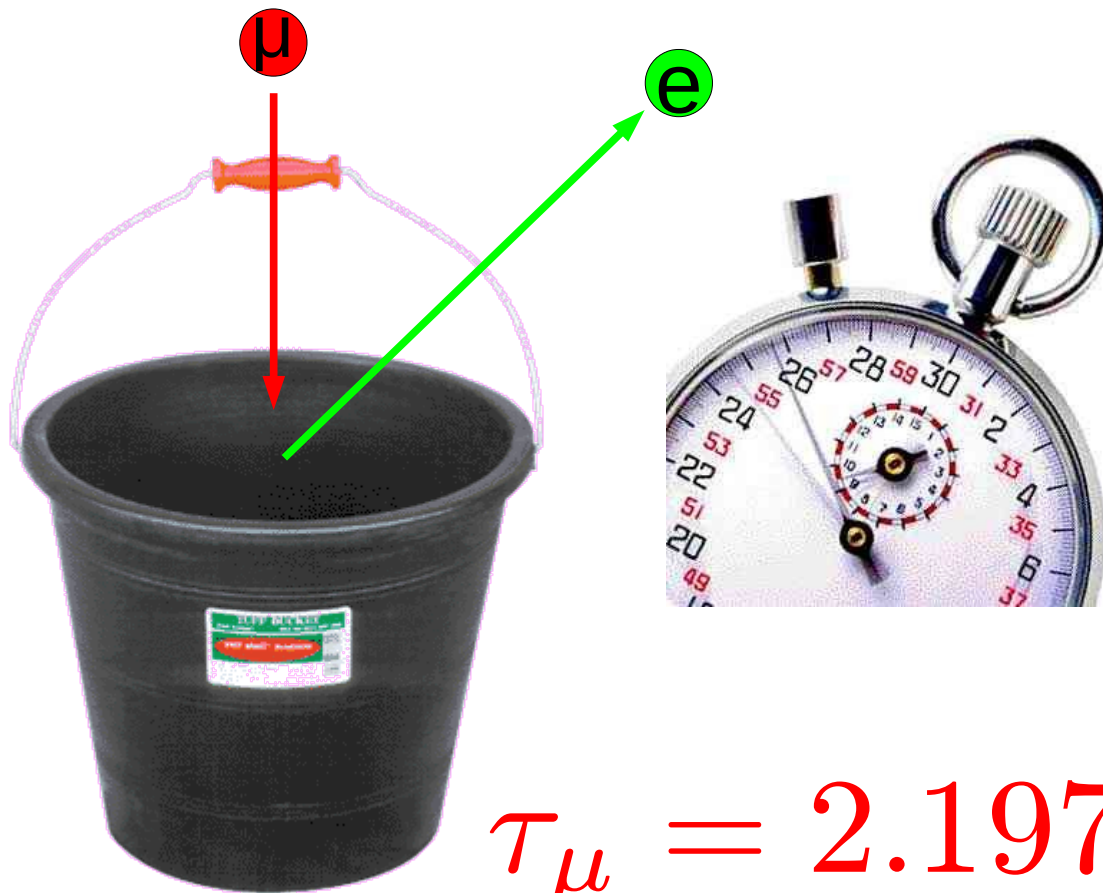
One-at-a-time



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One-at-a-time

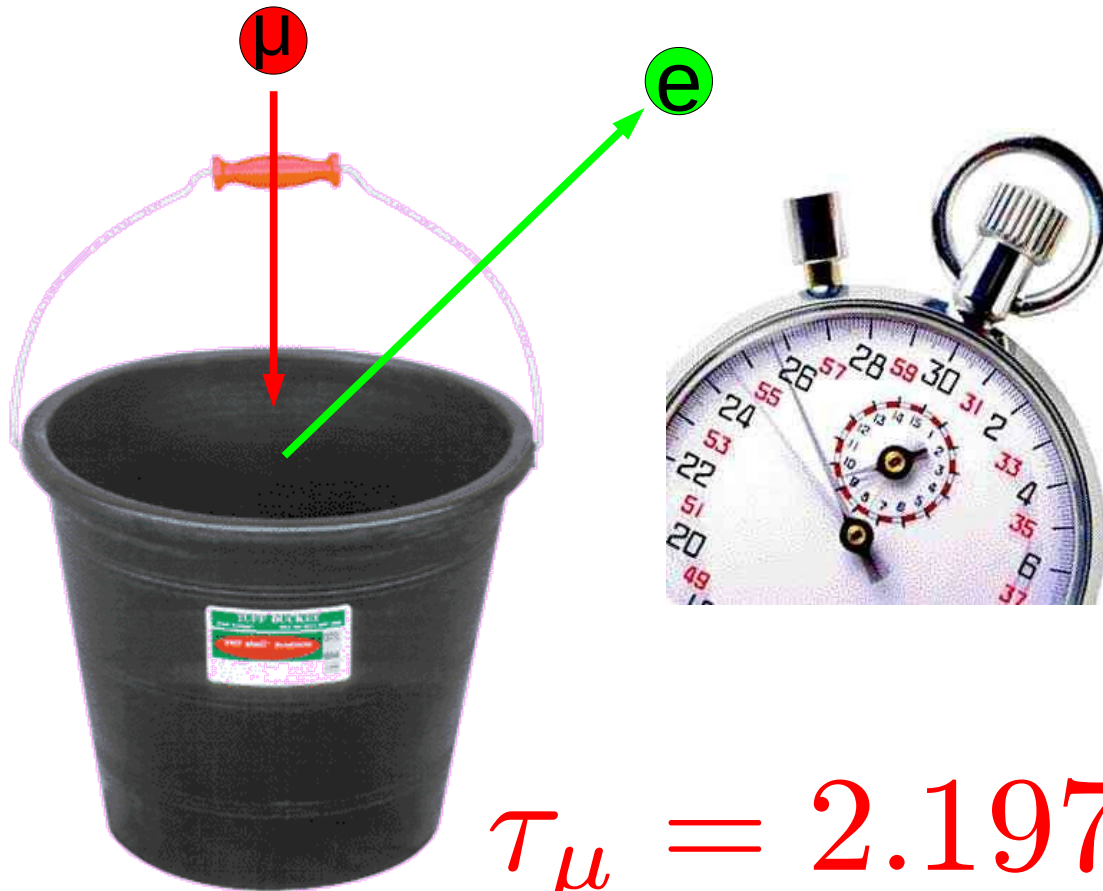


$$\tau_{\mu} = 2.197 \mu\text{s}$$

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One-at-a-time

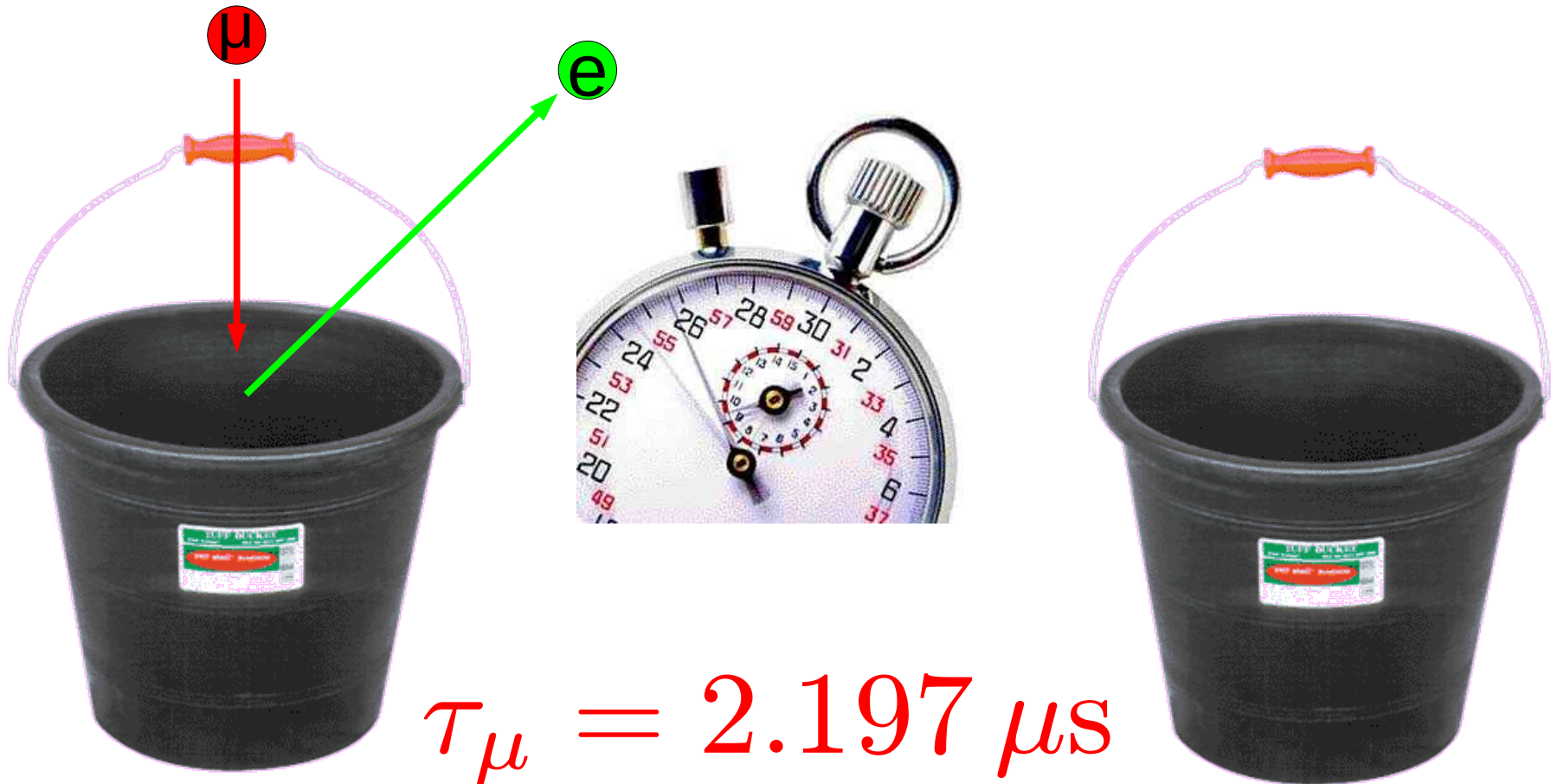
Many-at-once



How do you measure the muon lifetime?

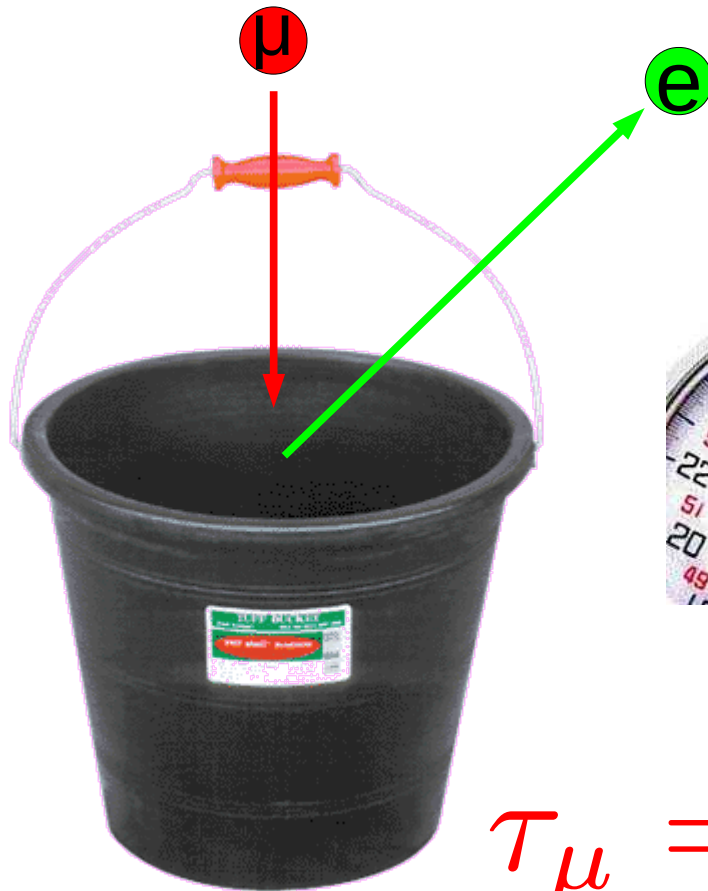
One-at-a-time

Many-at-once



How do you measure the muon lifetime?

One-at-a-time



Many-at-once

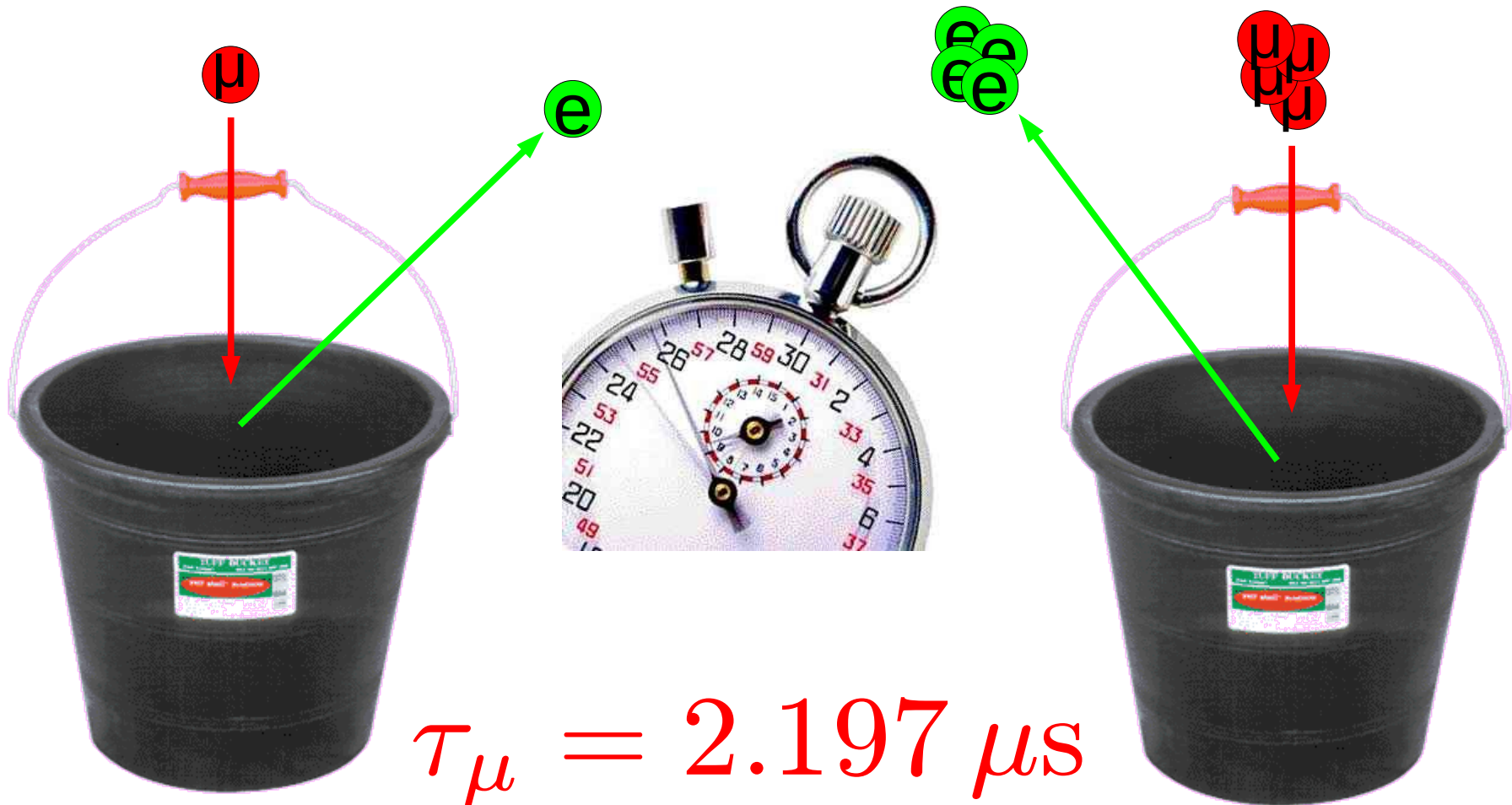


$$\tau_{\mu} = 2.197 \mu\text{s}$$

How do you measure the muon lifetime?

One-at-a-time

Many-at-once



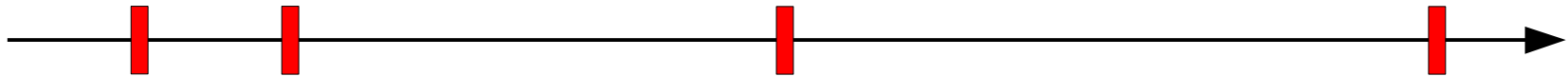
One-at-a-time

Can't really do one-at-a-time, the next best thing is a low rate, DC beam.

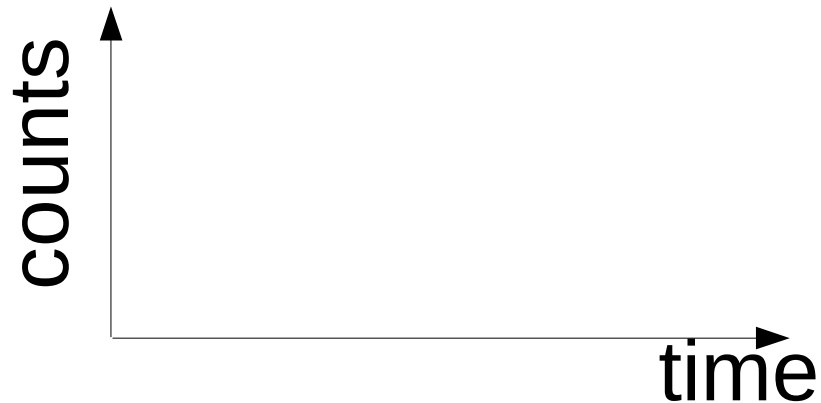
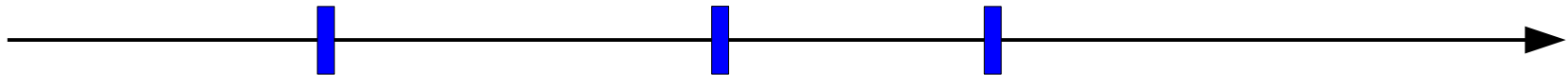
One-at-a-time

Can't really do one-at-a-time, the next best thing is a low rate, DC beam.

Muon timeline

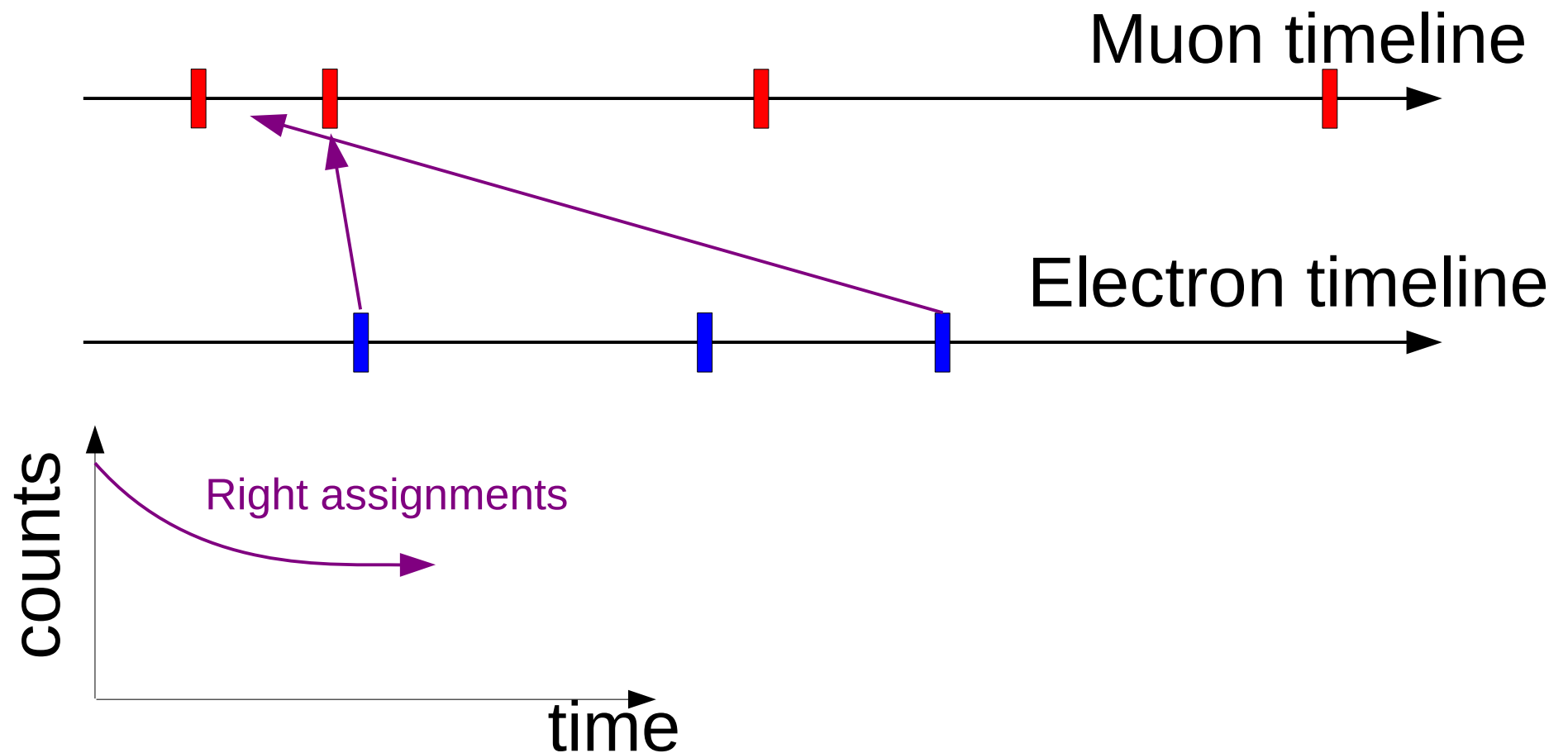


Electron timeline



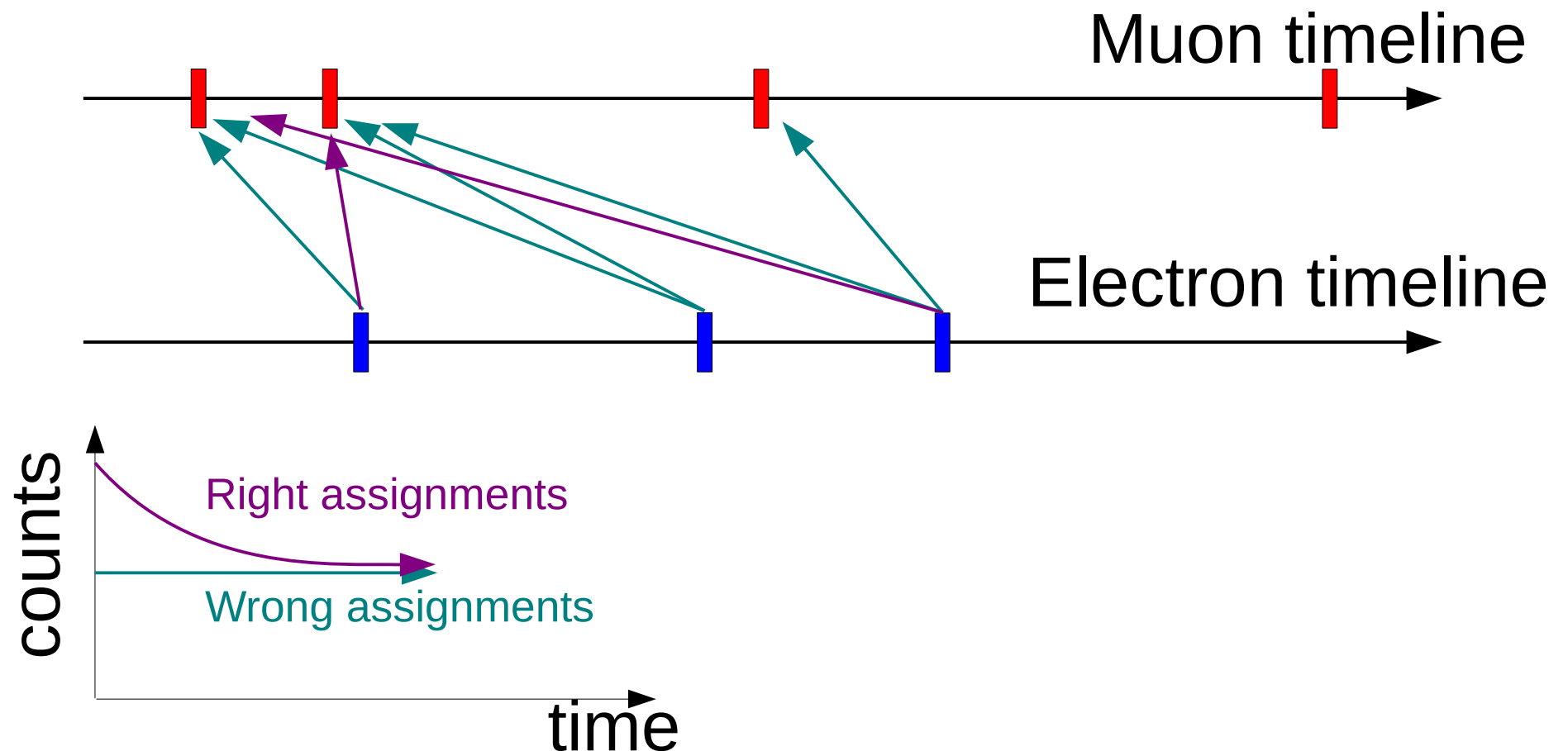
One-at-a-time

Can't really do one-at-a-time, the next best thing is a low rate, DC beam.



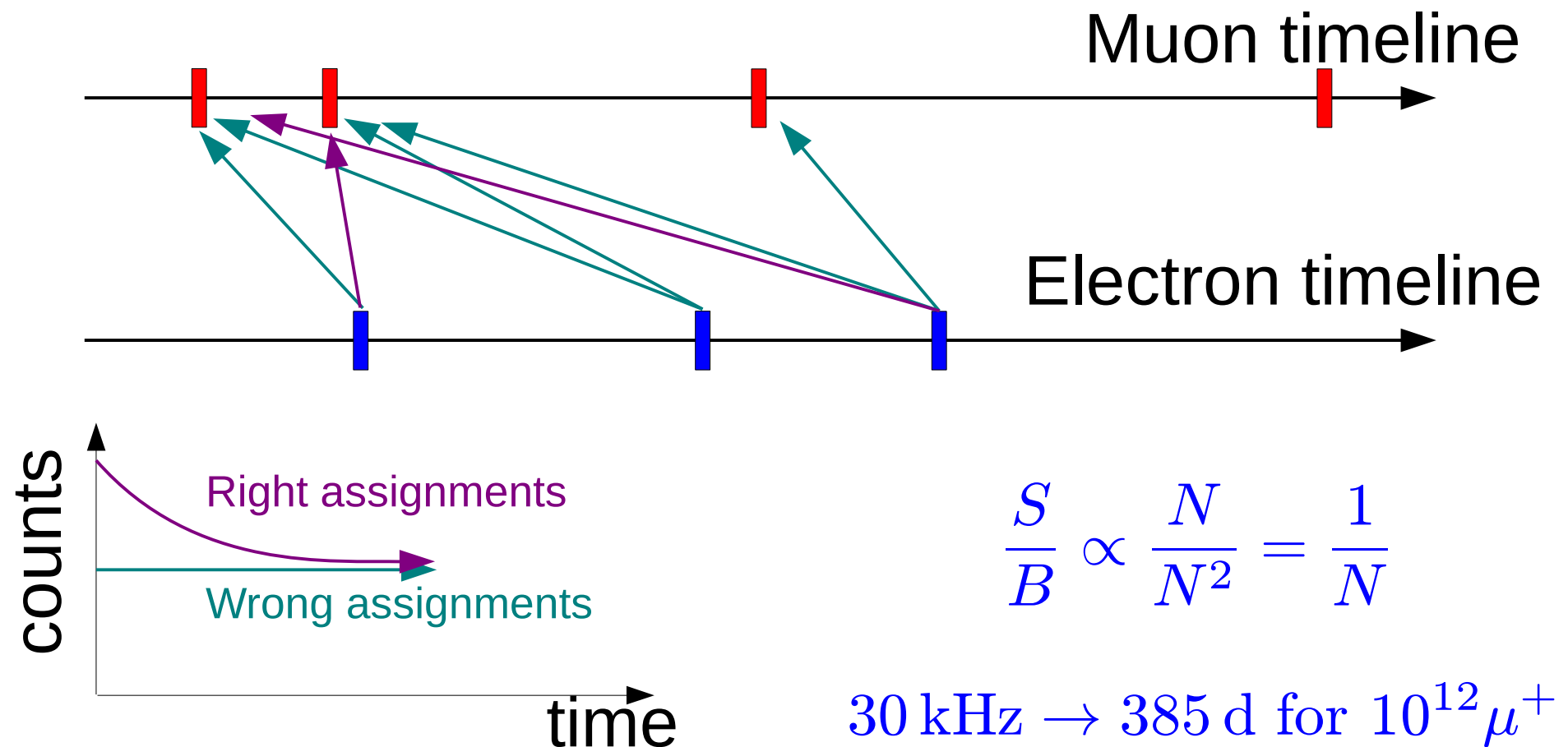
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One-at-a-time

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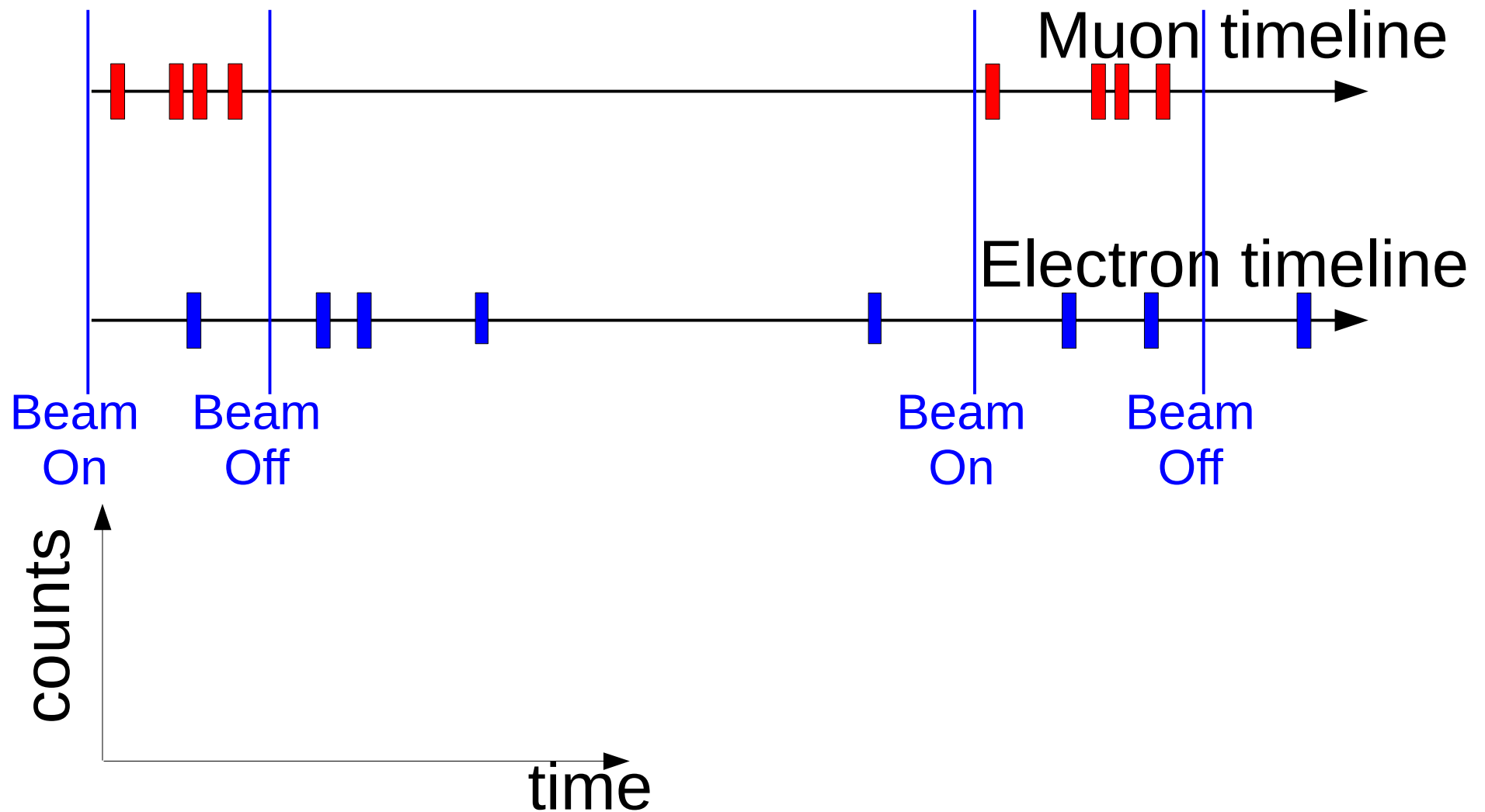


Many-at-once

Need time structured (AC) beam, not a continuous (DC) beam

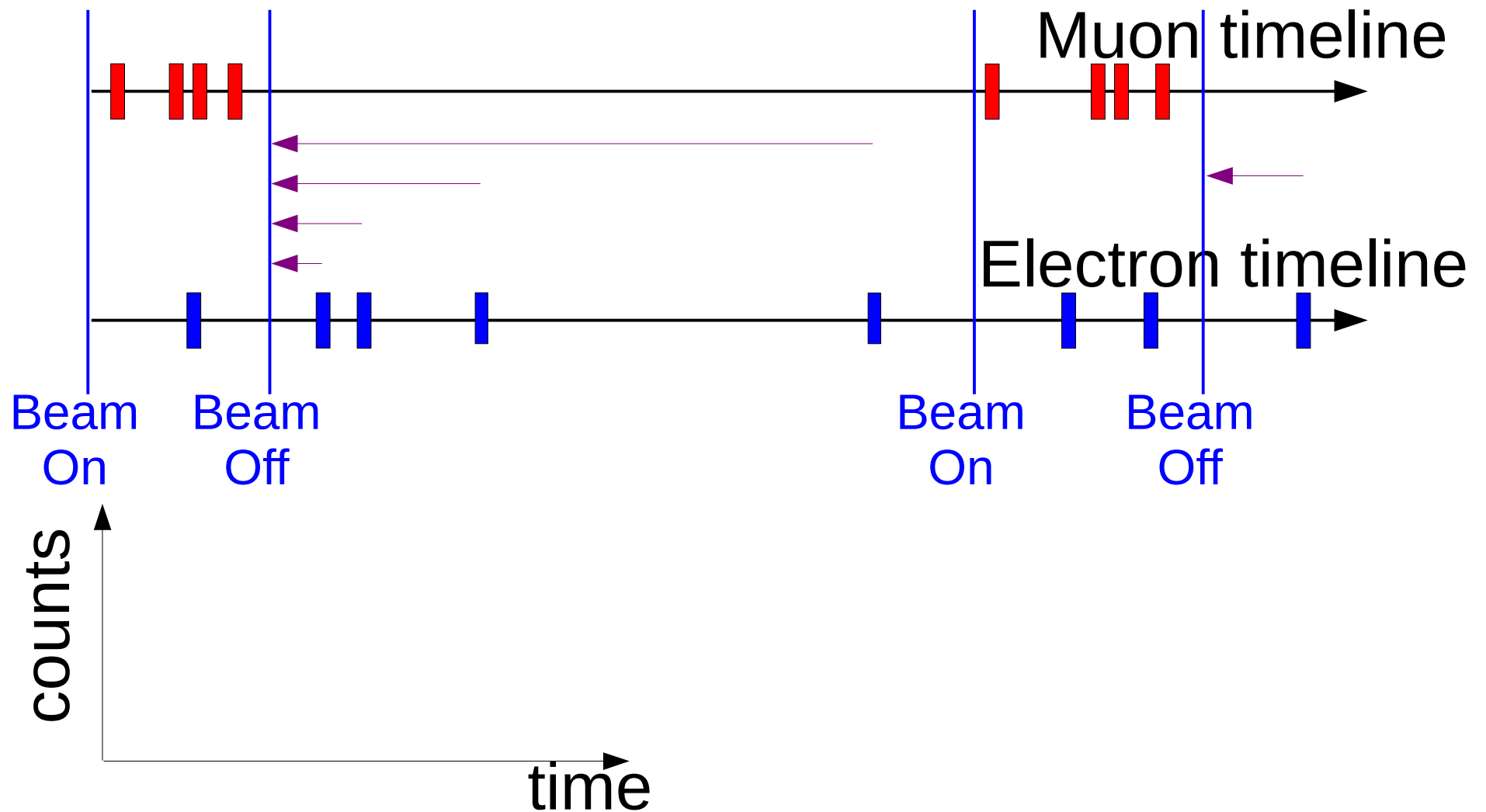
Many-at-once

Need time structured (AC) beam, not a continuous (DC) beam



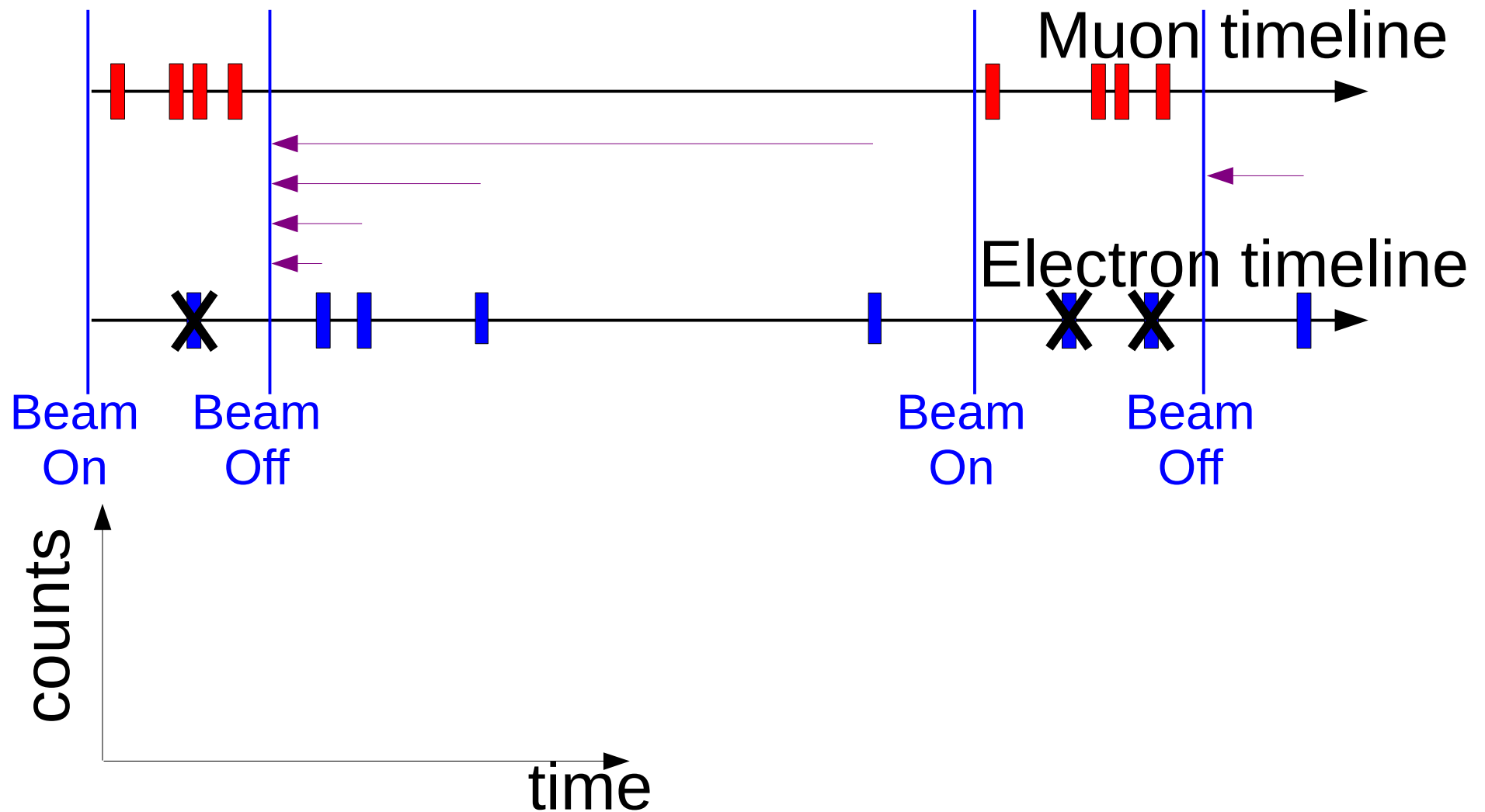
Many-at-once

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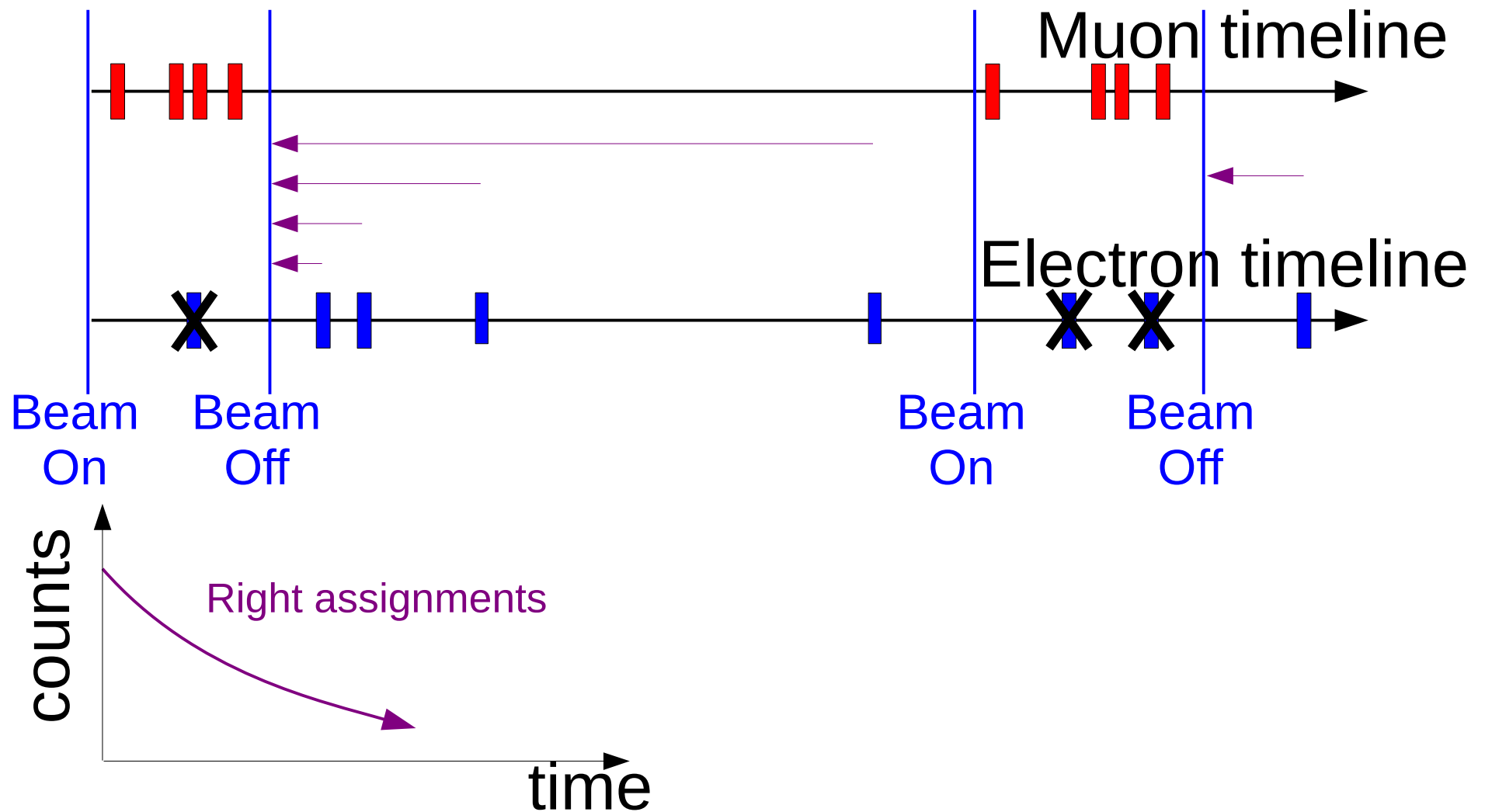
Many-at-once

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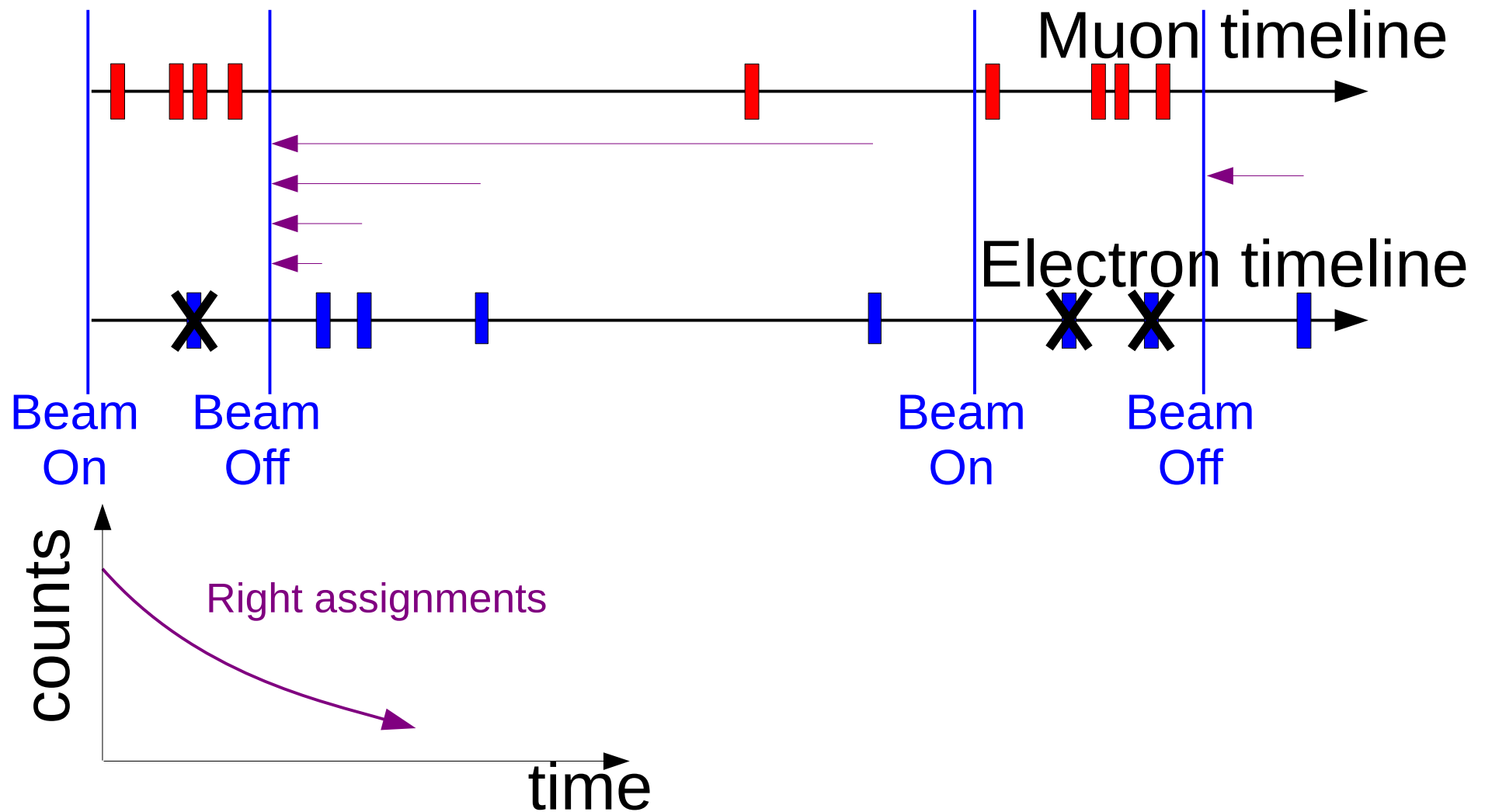
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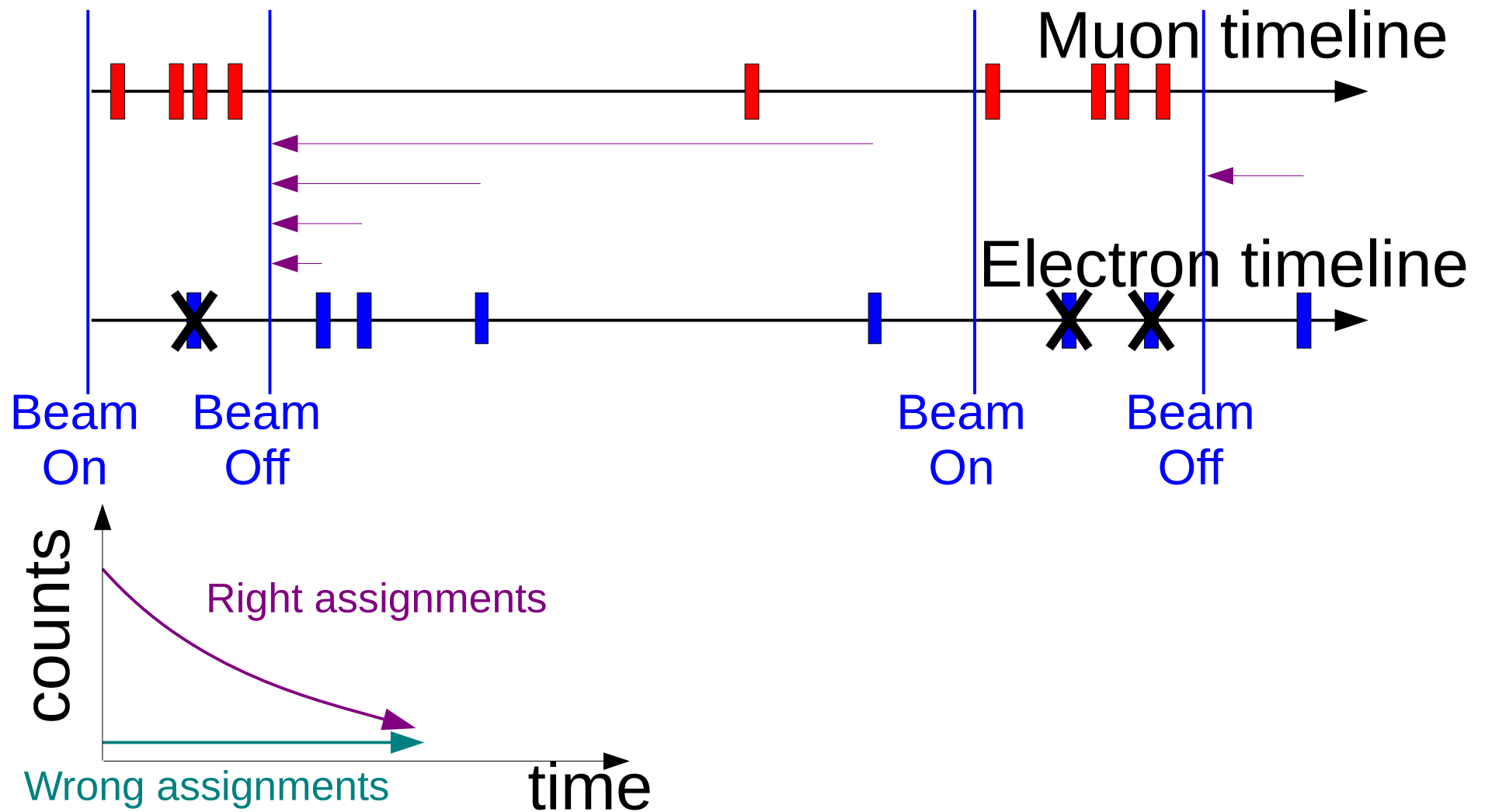
Many-at-once

Need time structured (AC) beam, not a continuous (DC) beam



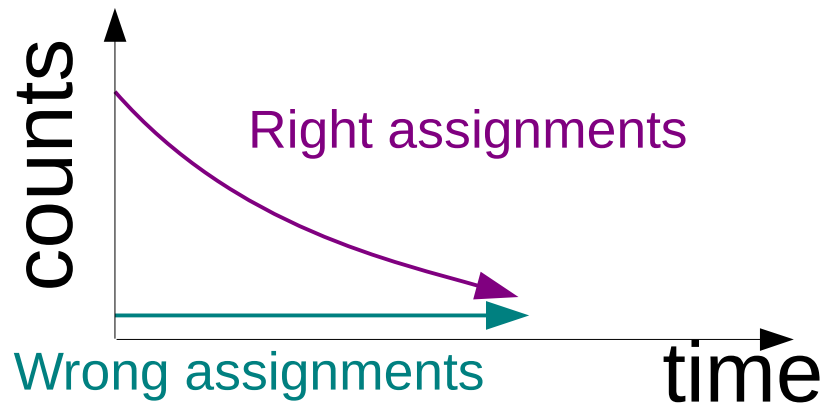
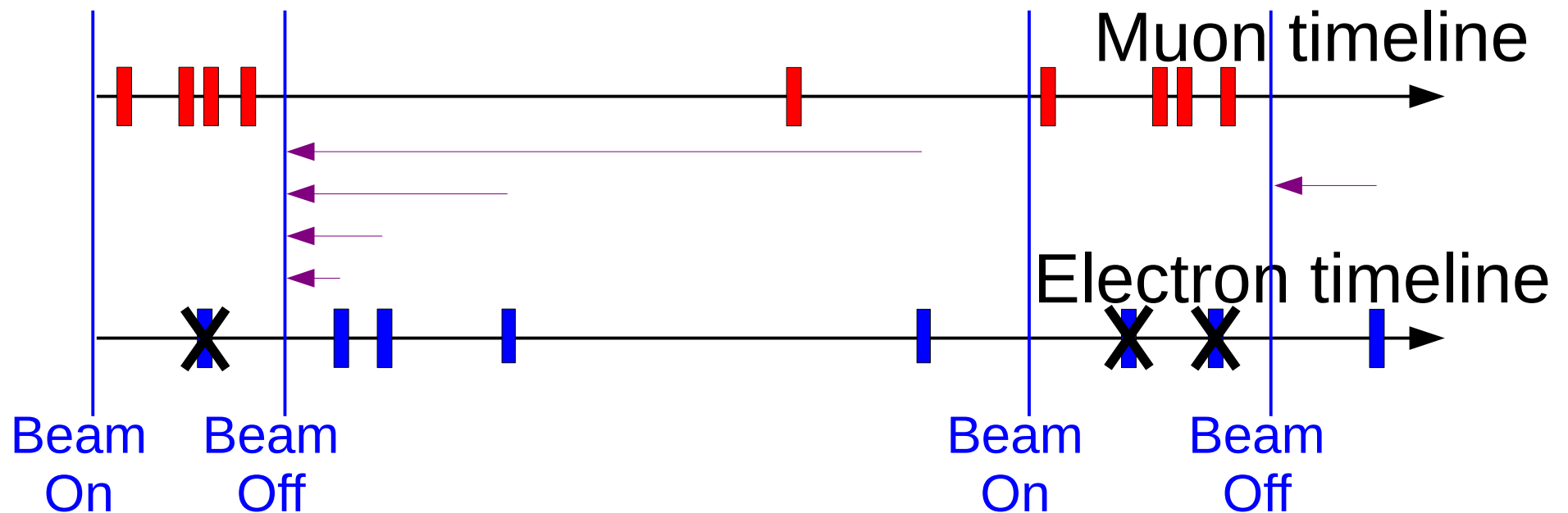
Many-at-once

Need time structured (AC) beam, not a continuous (DC) beam



Many-at-once

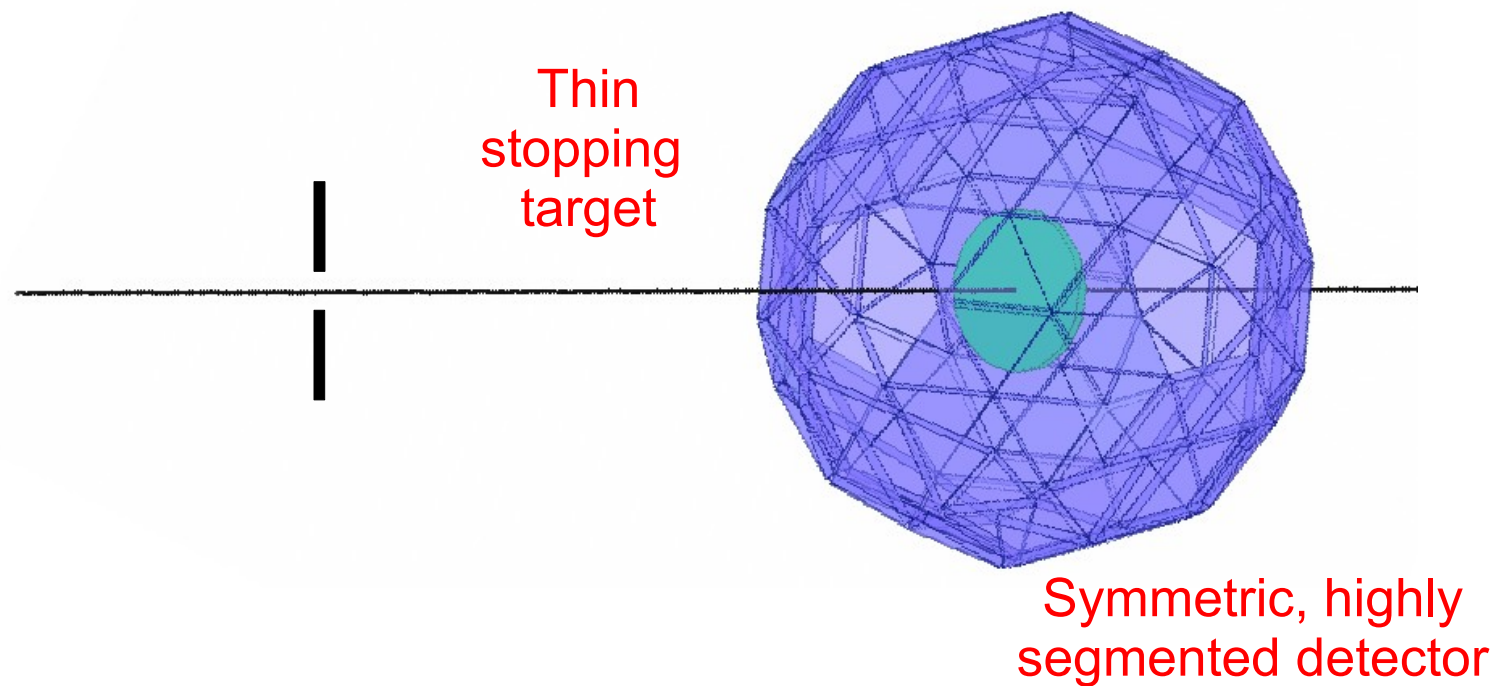
Need time structured (AC) beam, not a continuous (DC) beam



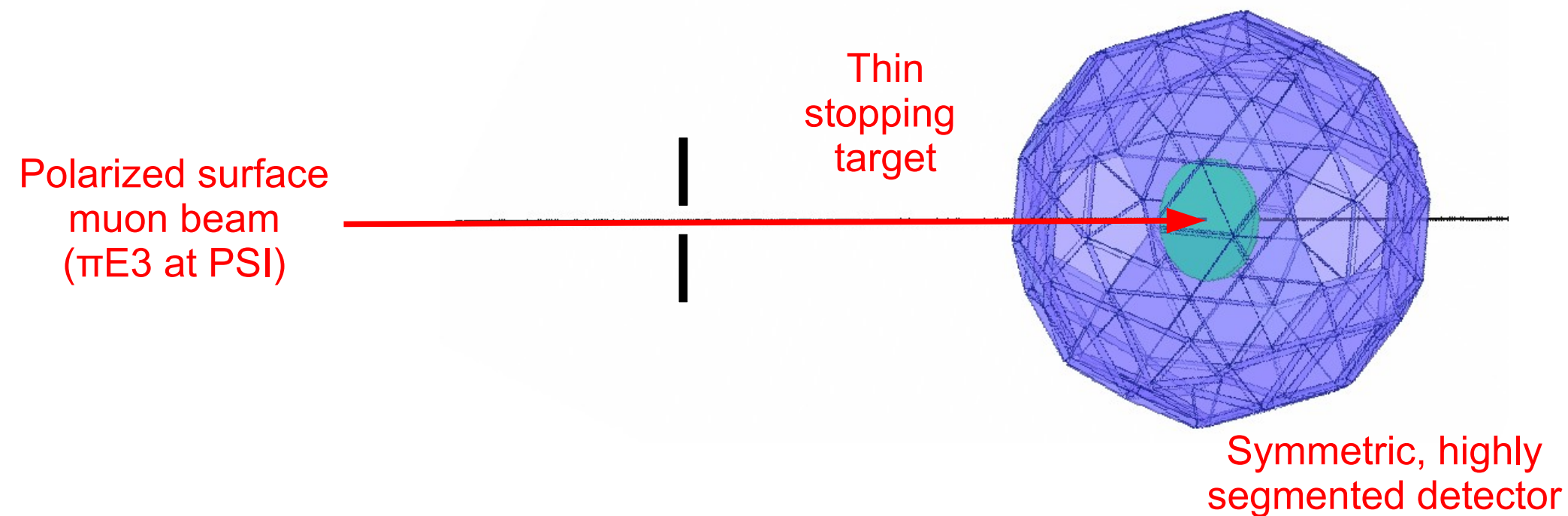
Much higher rates, but much harder experiment R&D and construction

We run in the “many-at-once” mode

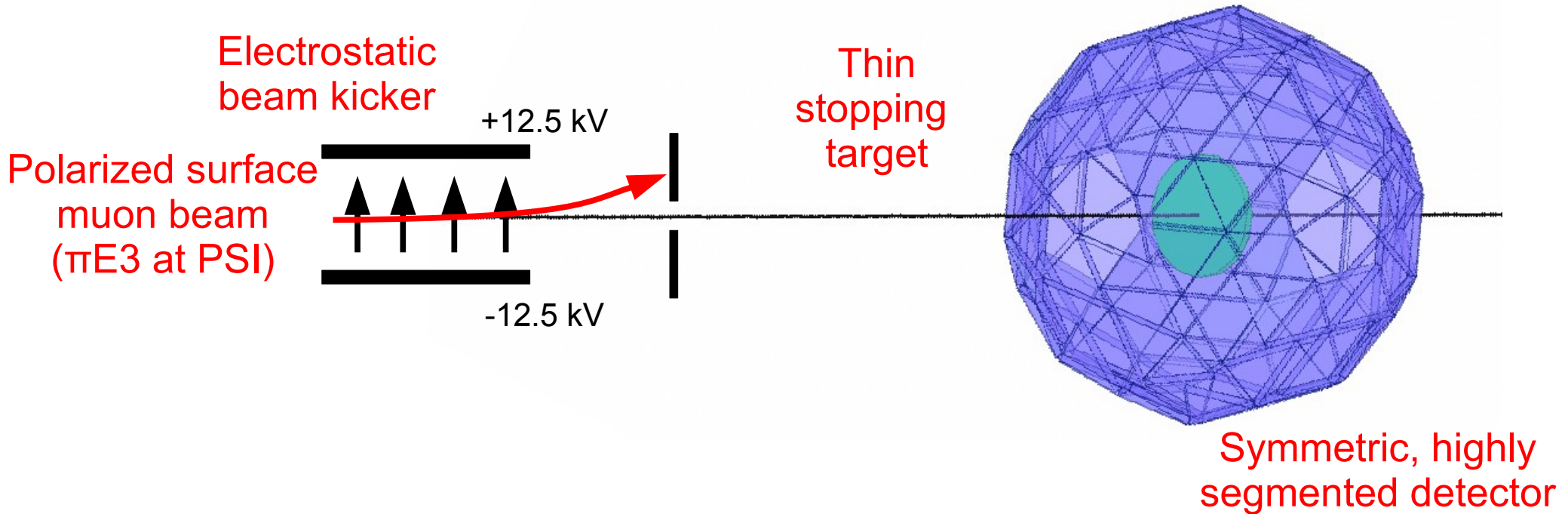
We run in the “many-at-once” mode



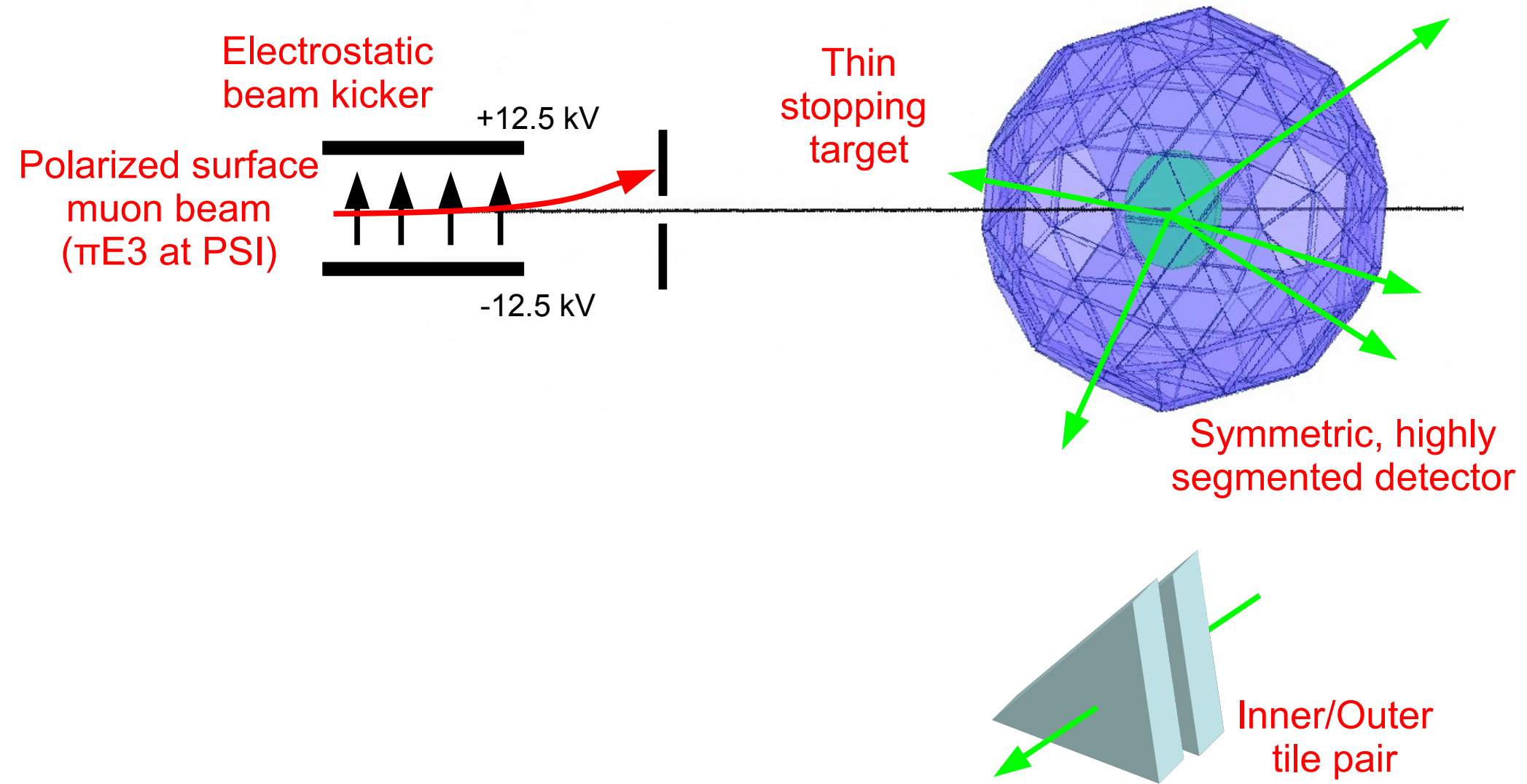
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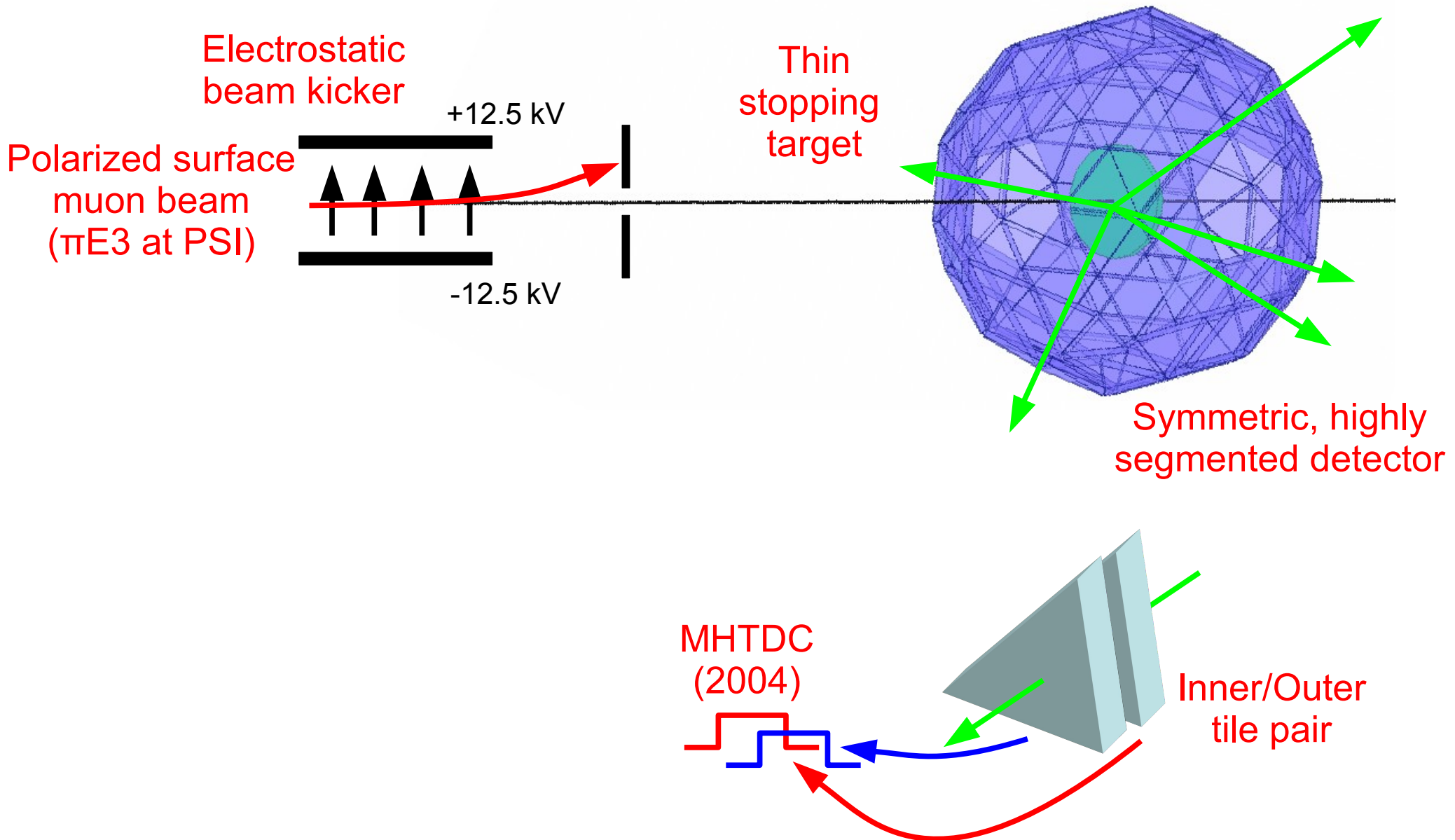
We run in the “many-at-once” mode



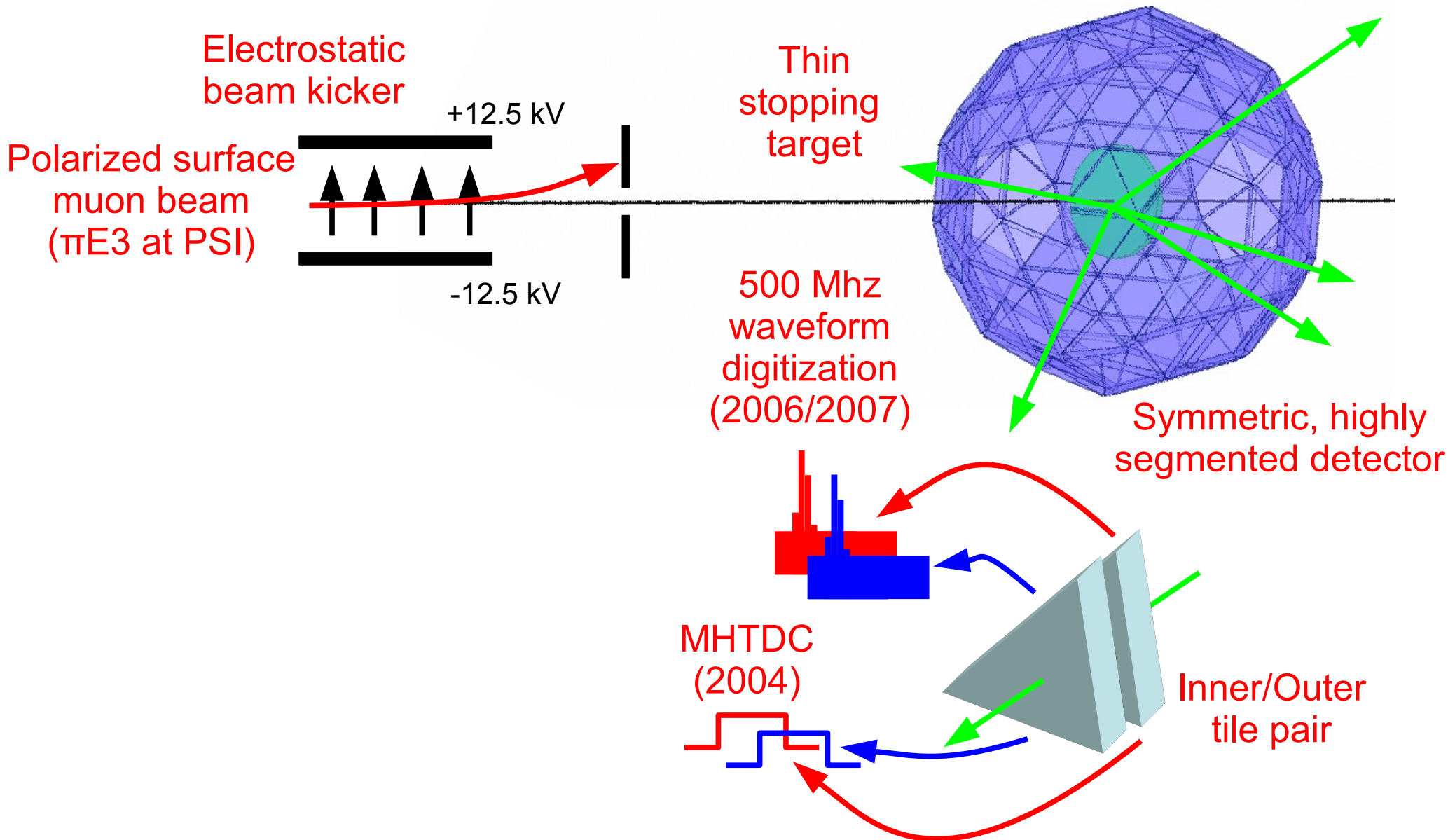
We run in the “many-at-once” mode



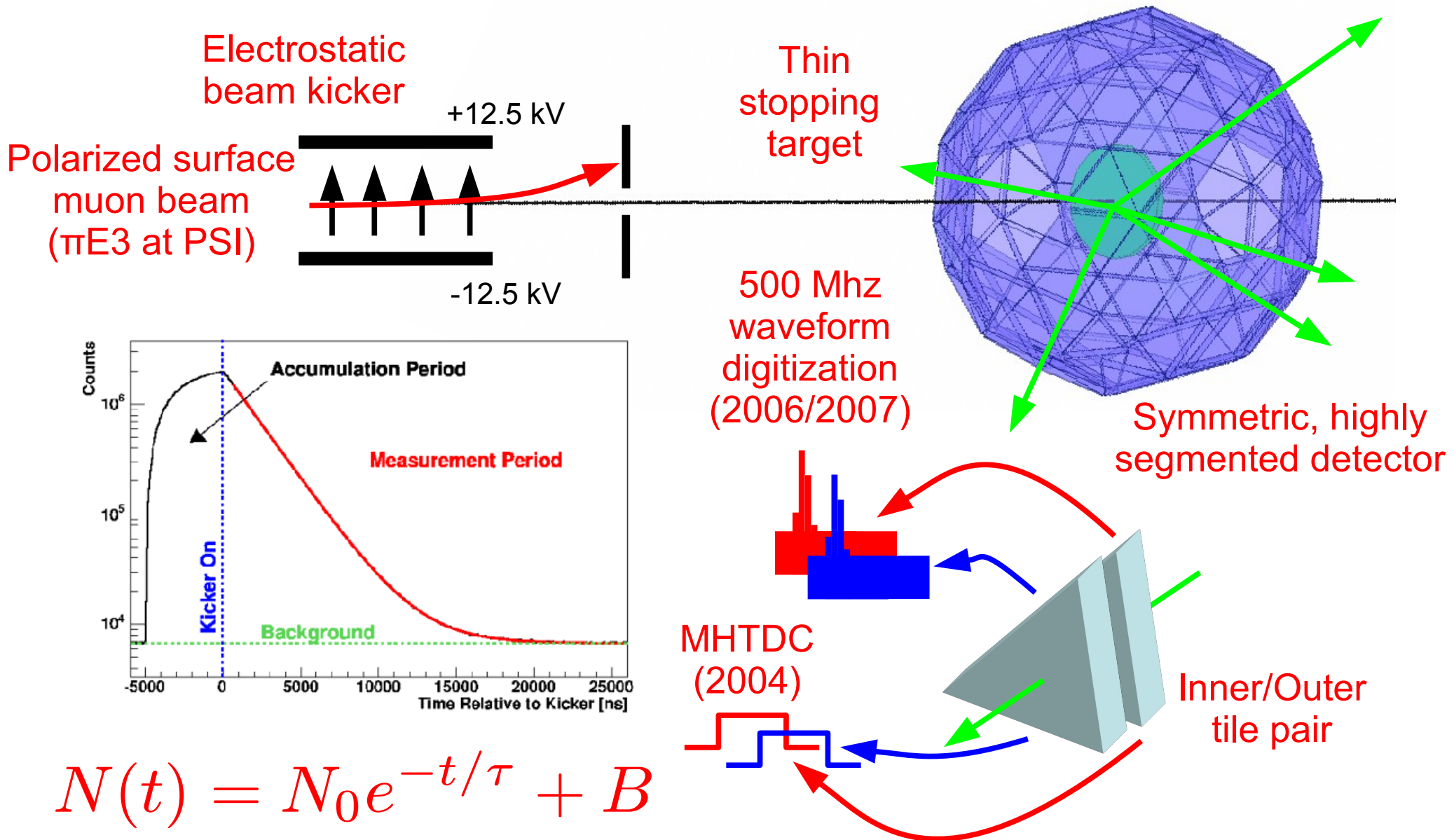
We run in the “many-at-once” mode



We run in the “many-at-once” mode



We run in the “many-at-once” mode



Time-dependent systematics are the core concern for a 10^{12} data set

Early-to-late changes, for instance:

Instrumental issues

PMT gains

Discriminator threshold walk

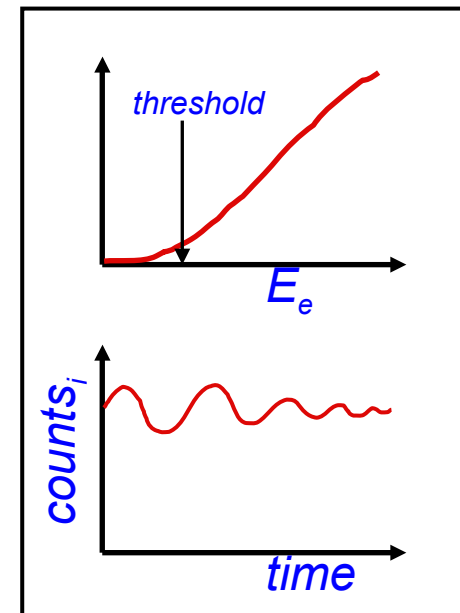
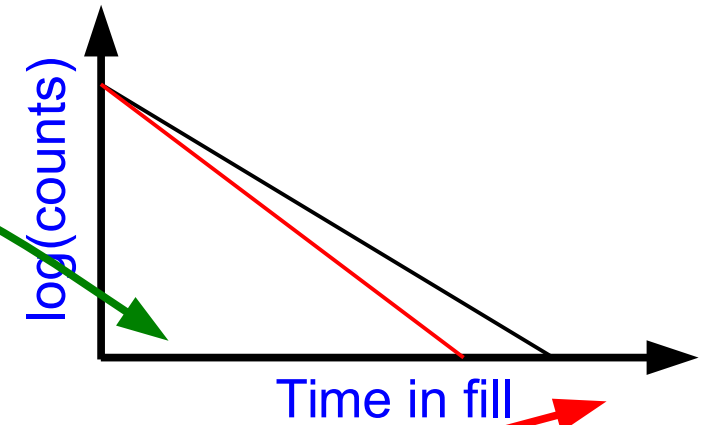
Kicker voltage sag

Pileup

Physics issues

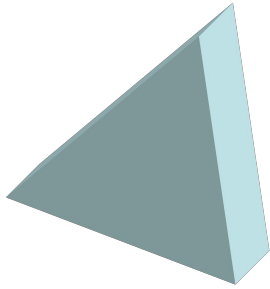
Polarization precession

Longitudinal relaxation



The pileup spectrum can be constructed directly from the data

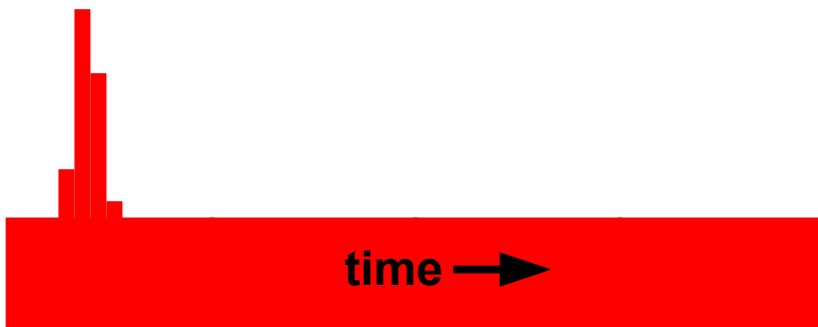
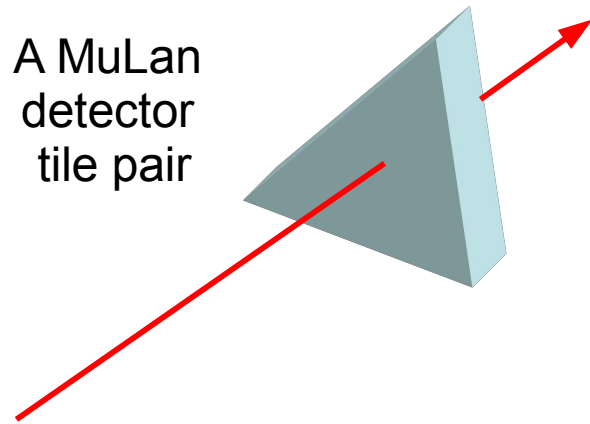
A MuLan
detector
tile pair



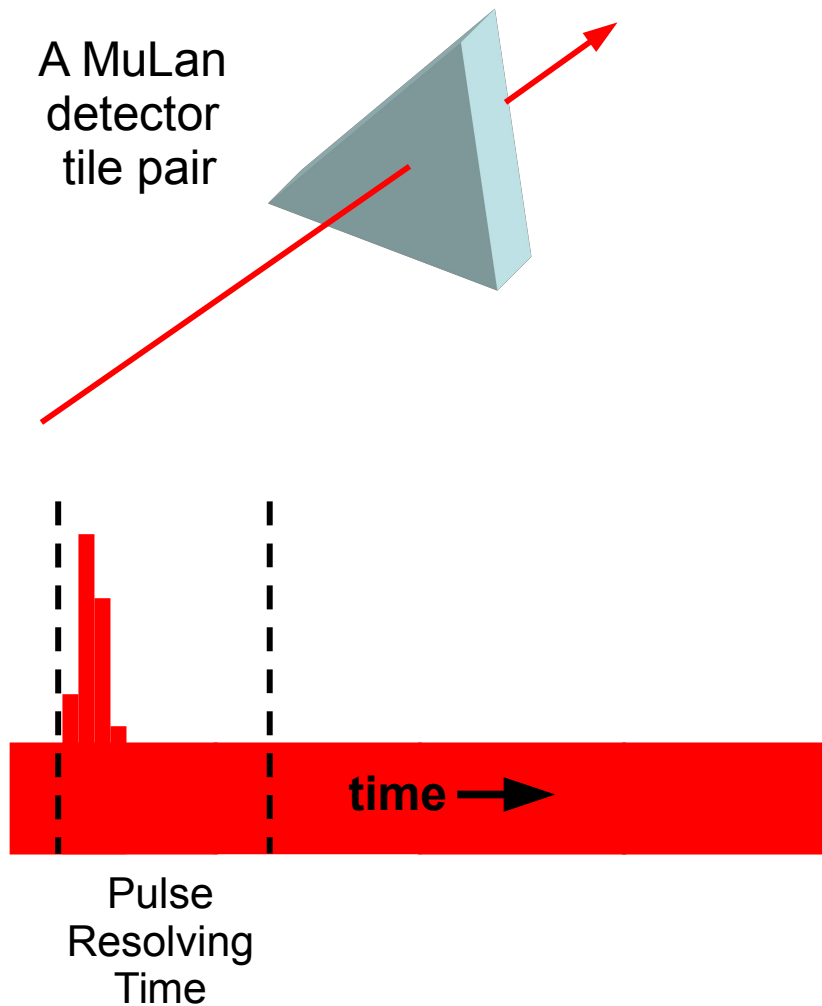
time →



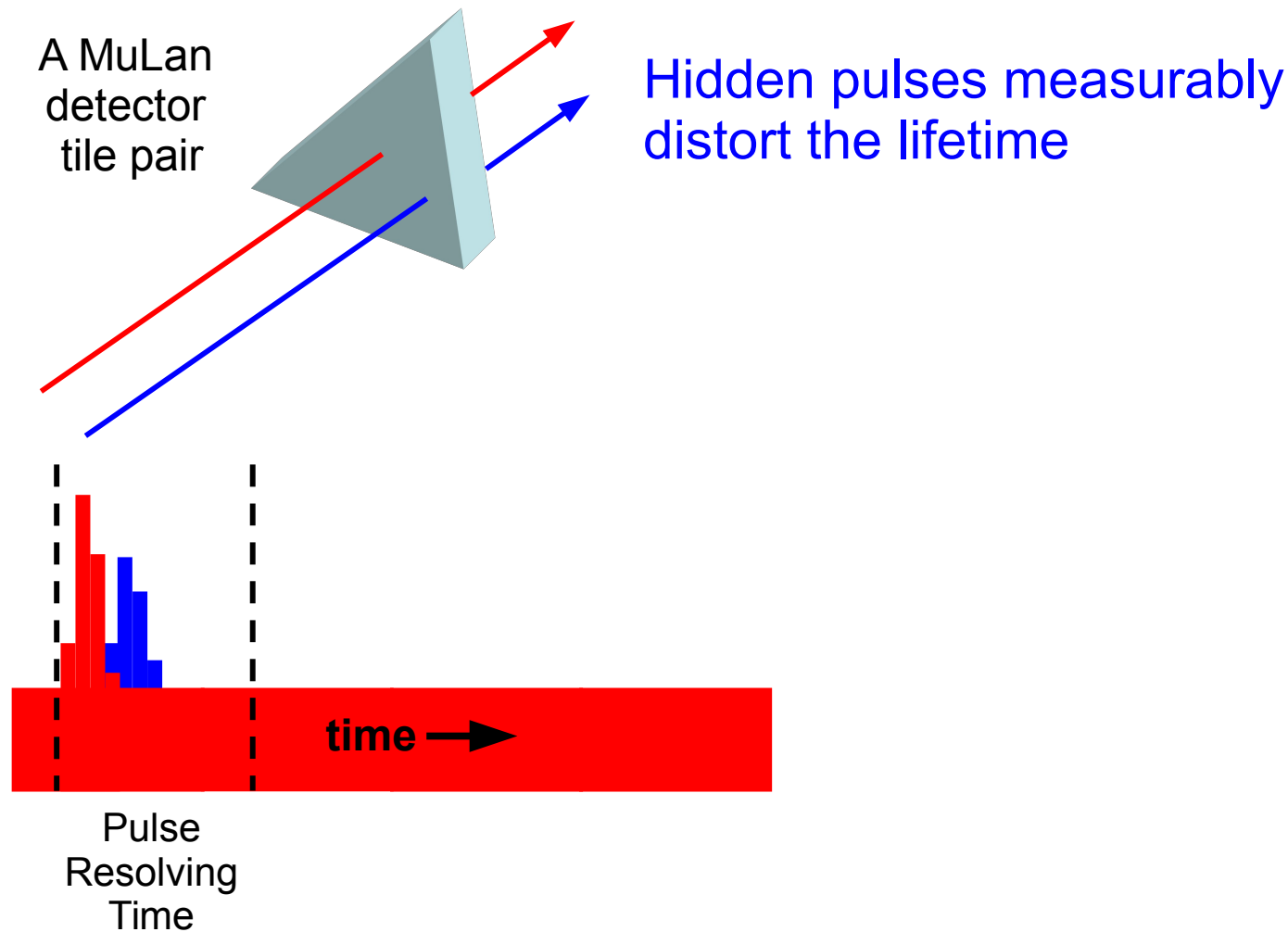
The pileup spectrum can be constructed directly from the data



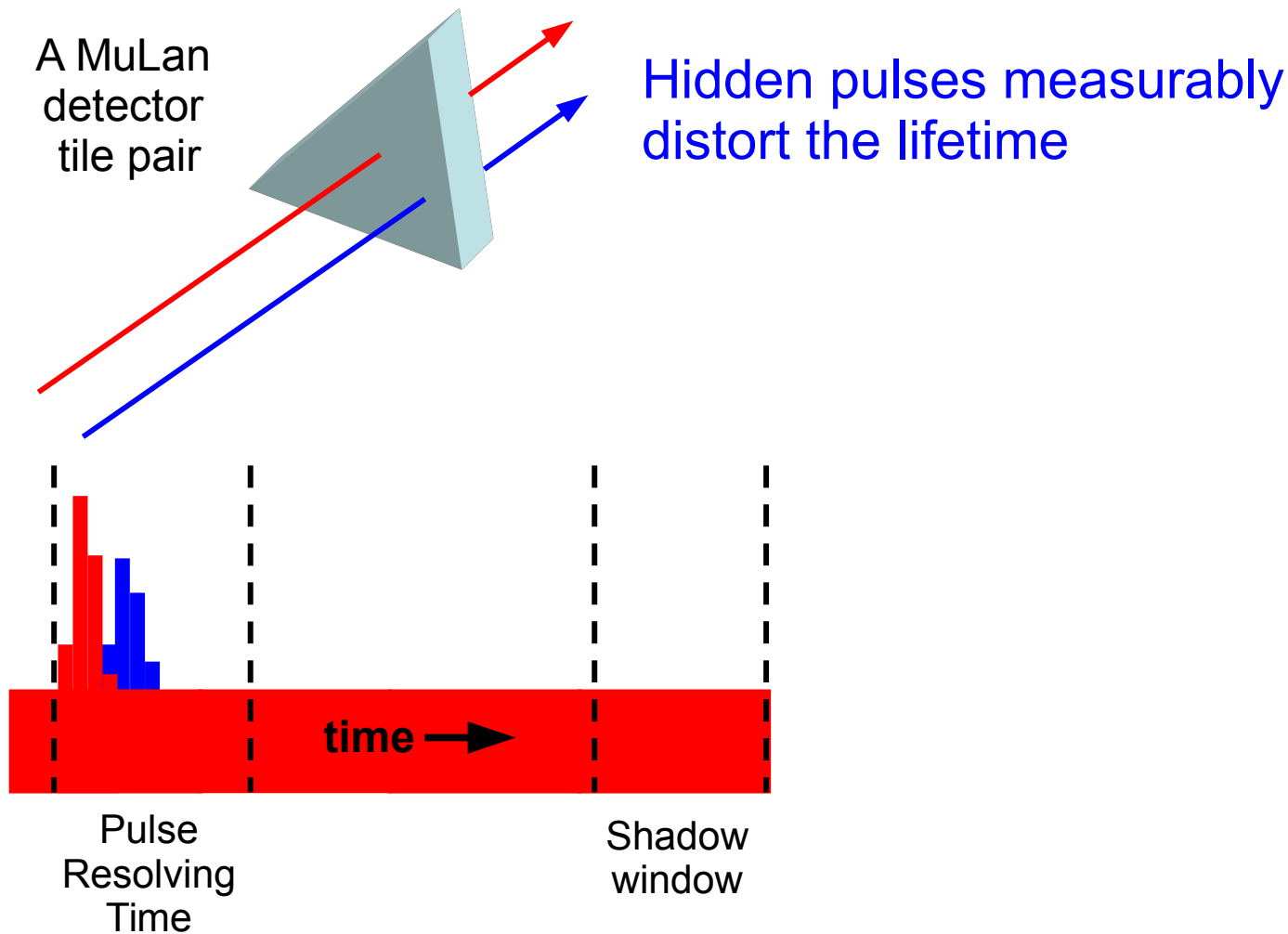
The pileup spectrum can be constructed directly from the data



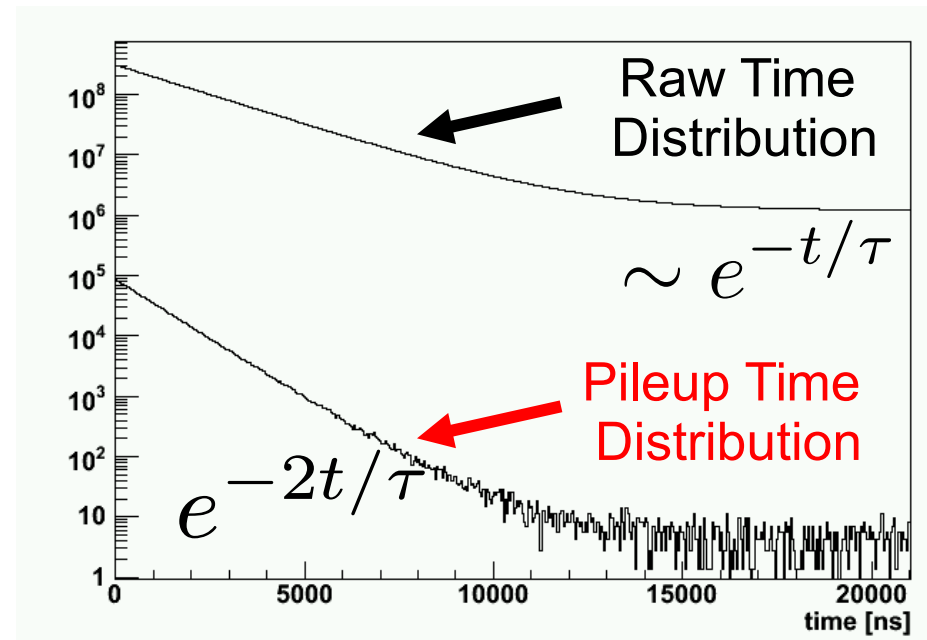
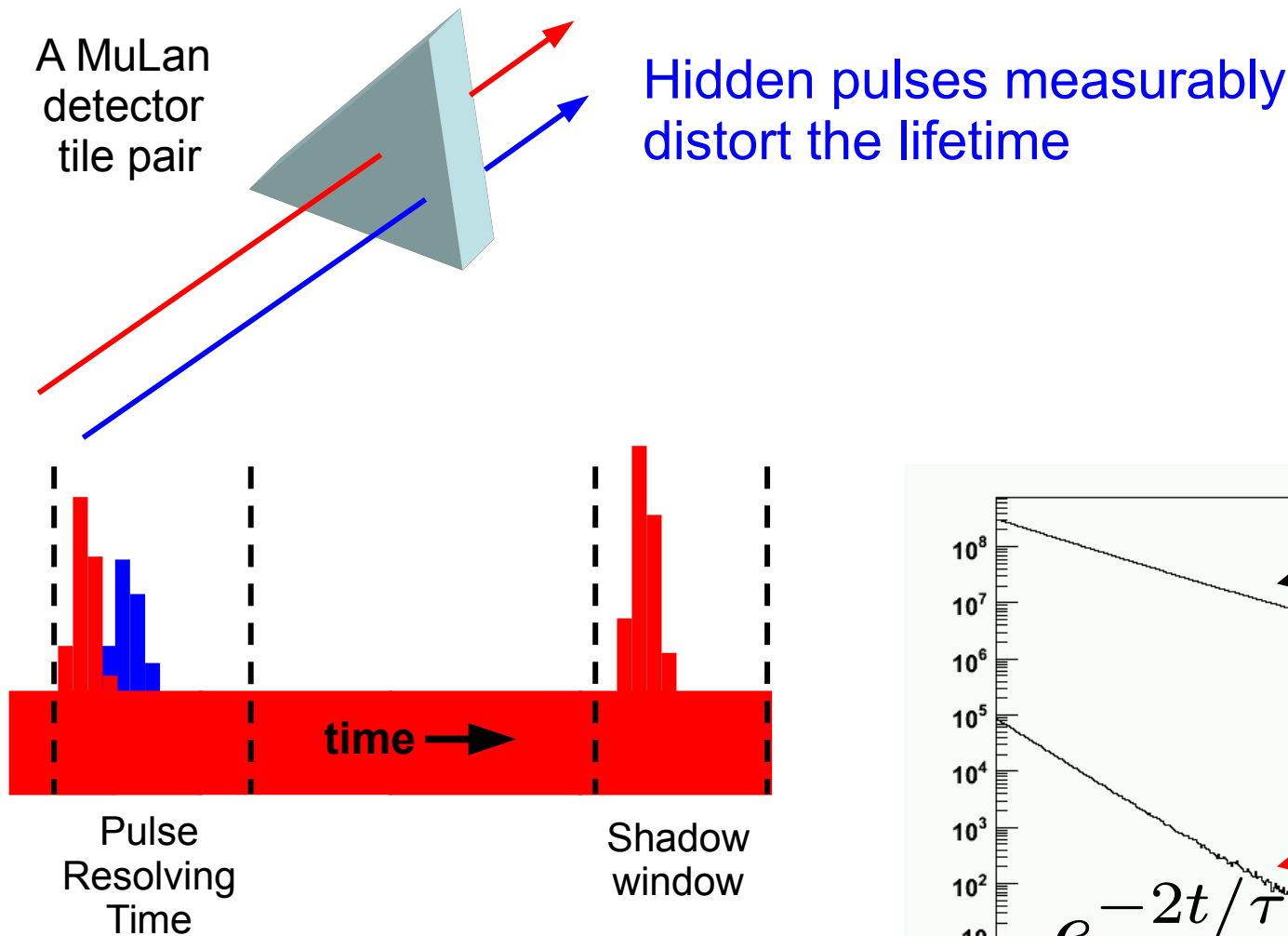
The pileup spectrum can be constructed directly from the data



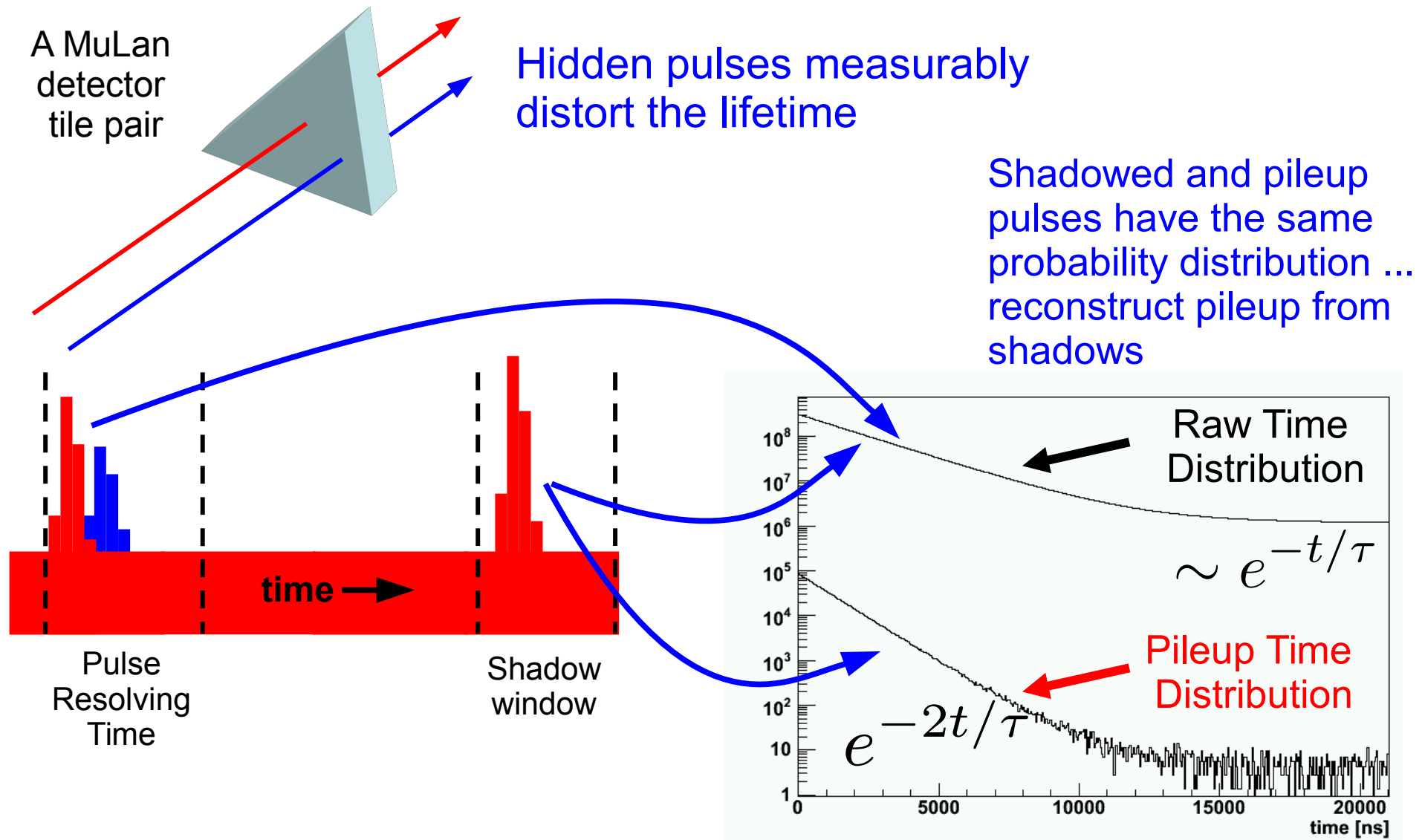
The pileup spectrum can be constructed directly from the data



The pileup spectrum can be constructed directly from the data

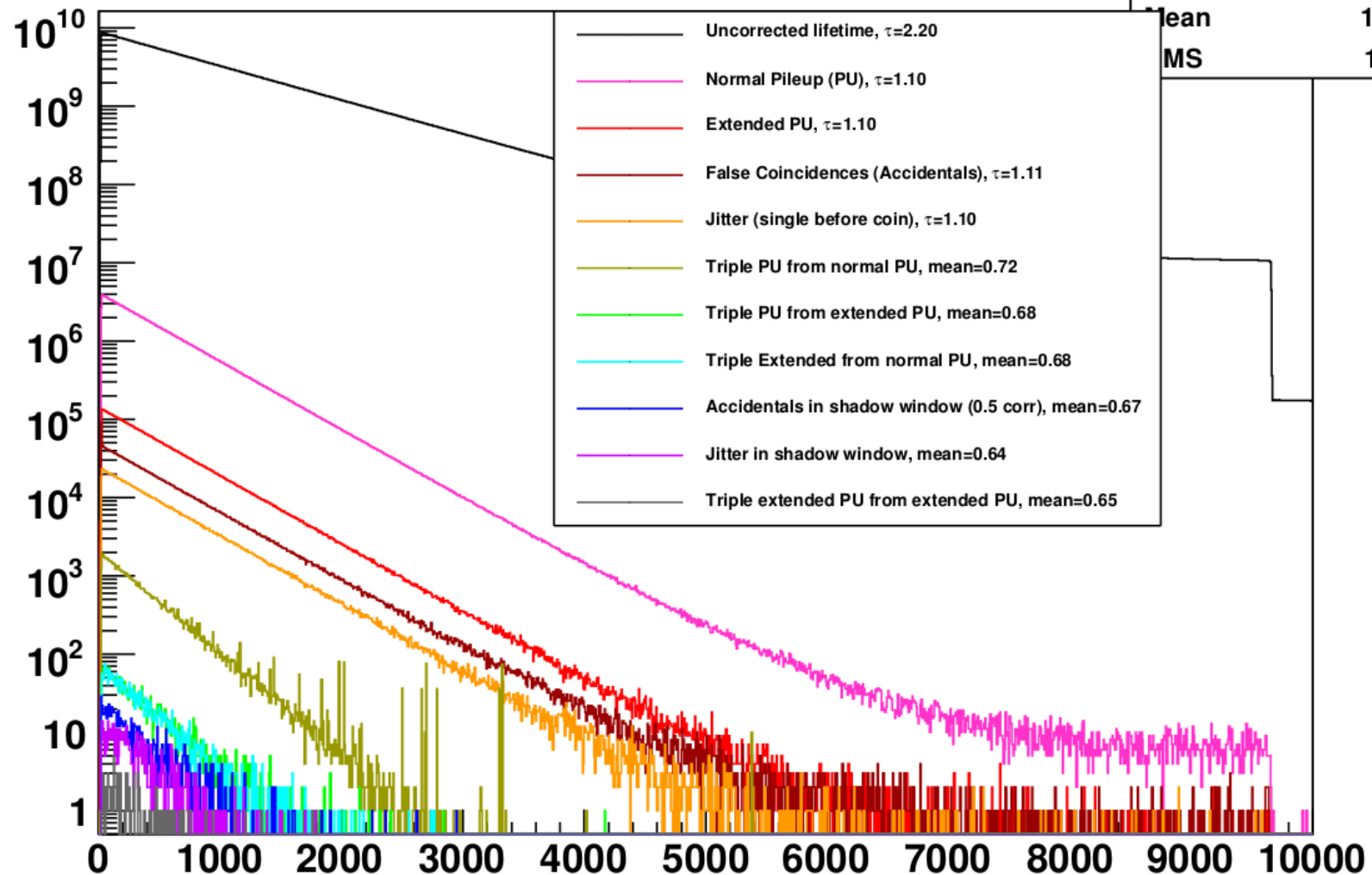


The pileup spectrum can be constructed directly from the data

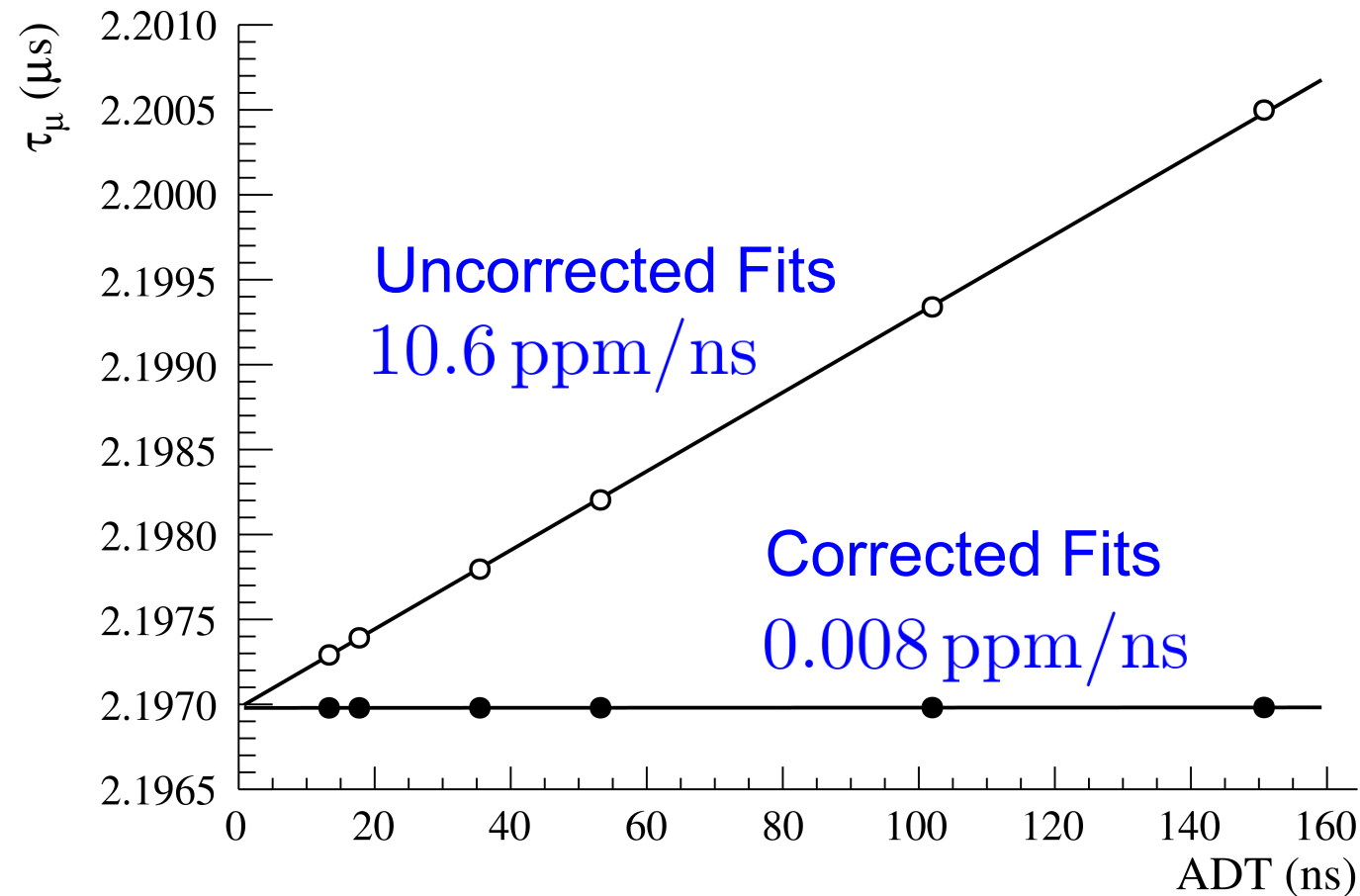


For a 1ppm measurement, we have to go well beyond the single event pileup spectrum

lifetimeLast ADT=5.00, CW=5.00



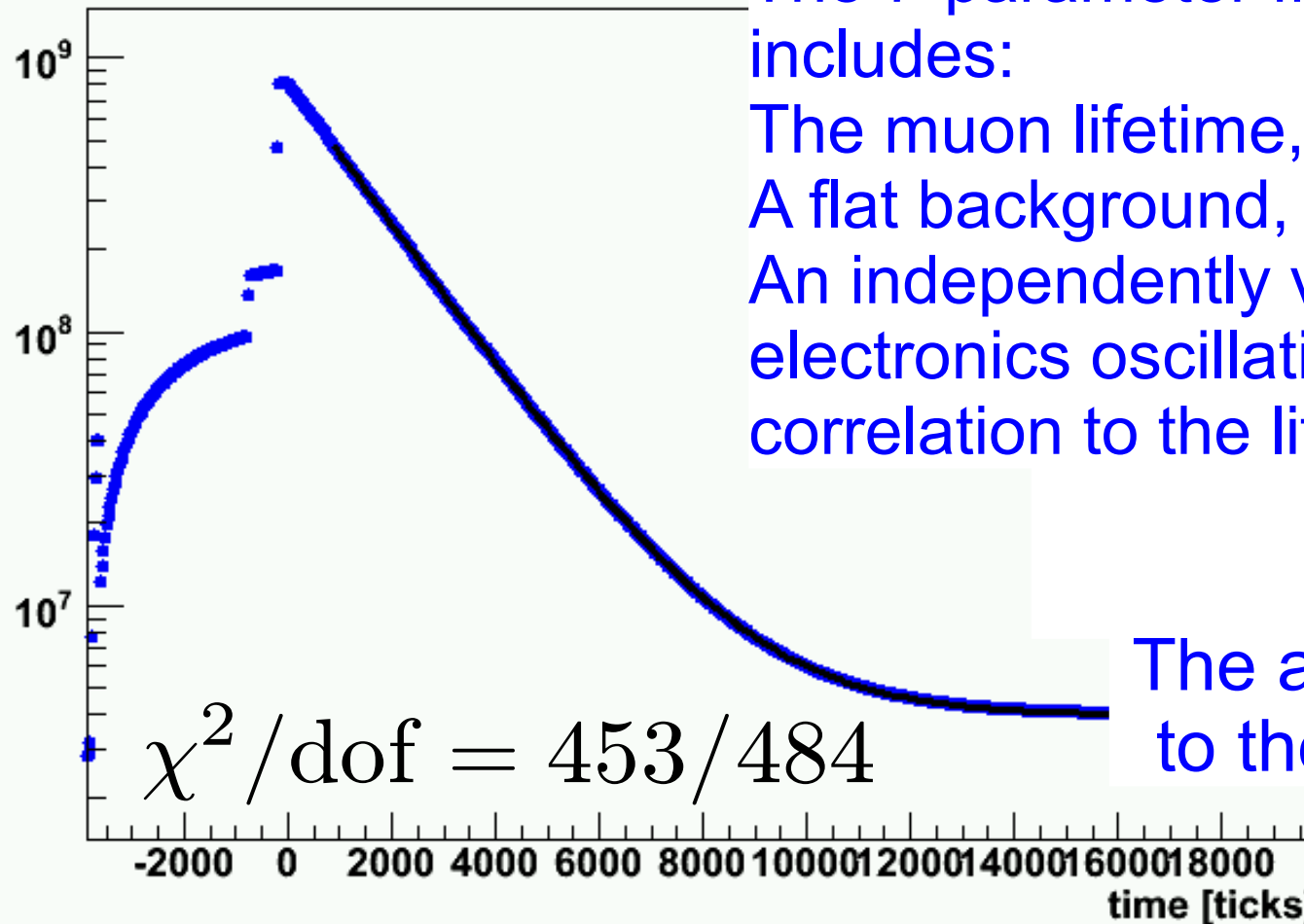
How well does the shadow method correct pileup?



In our final result, we extrapolate to zero resolution time, and apply a systematic to cover the very small residual effect

Analysis of our 2004 Physics run yielded a 11 ppm lifetime measurement

Lifetime Histogram



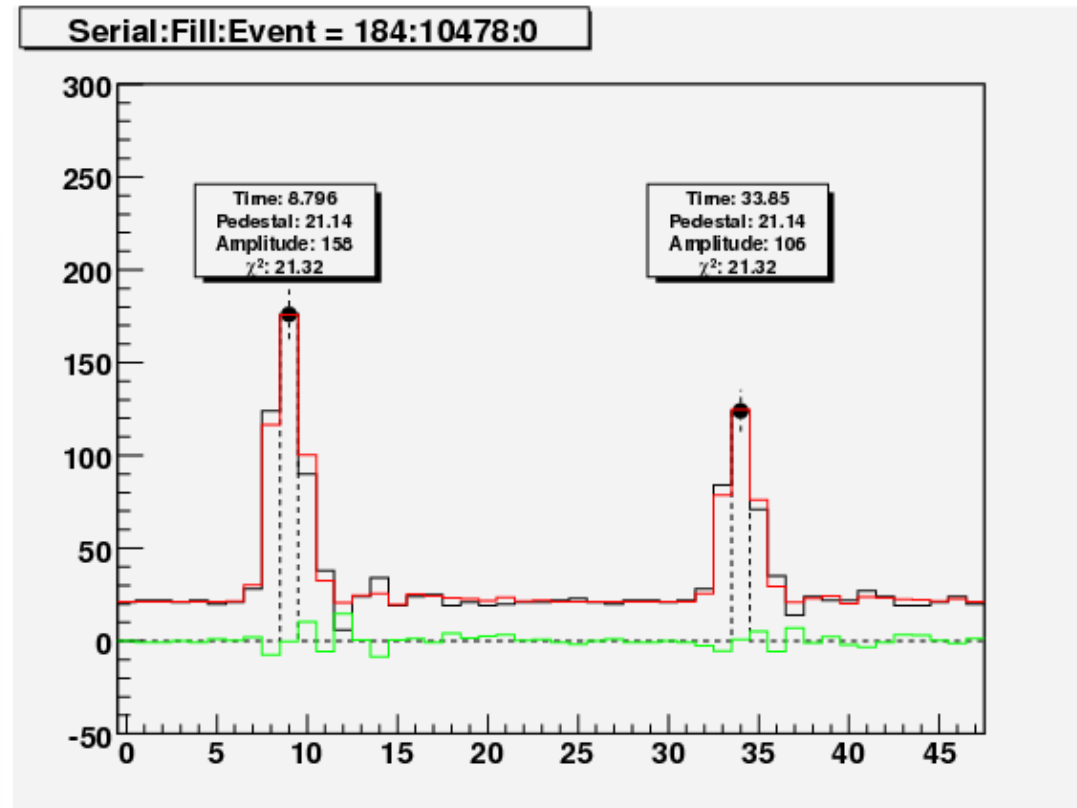
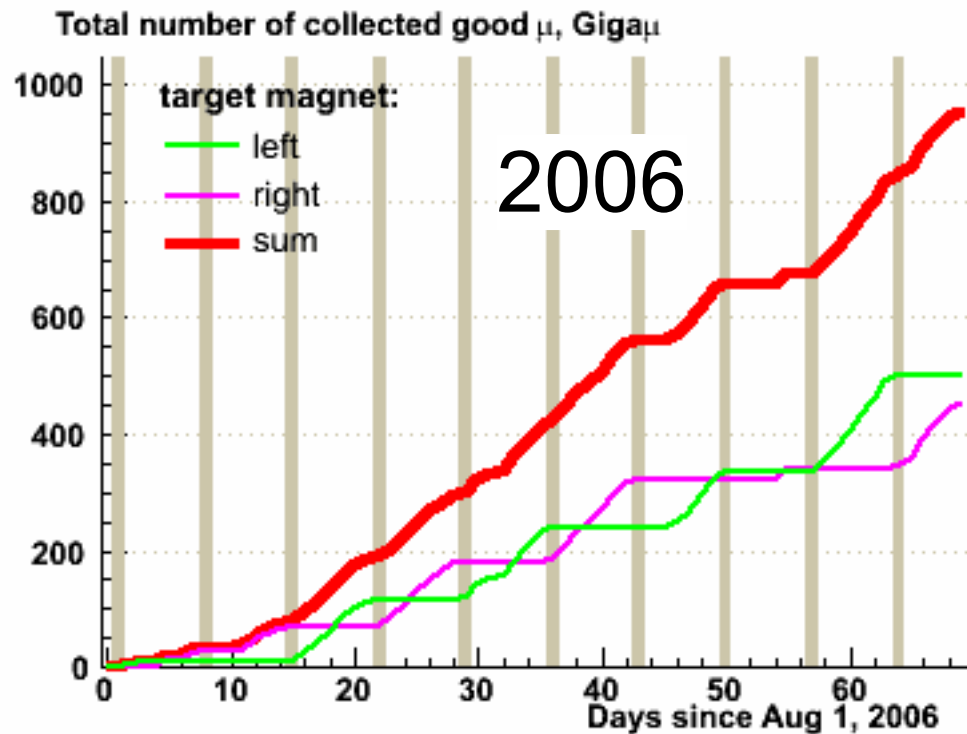
The 7-parameter fit function includes:

- The muon lifetime,
- A flat background, and
- An independently validated electronics oscillation (with low correlation to the lifetime)

The analyzers are blind to the clock frequency

We engaged in extensive analysis of our 2006 and 2007 data sets ...

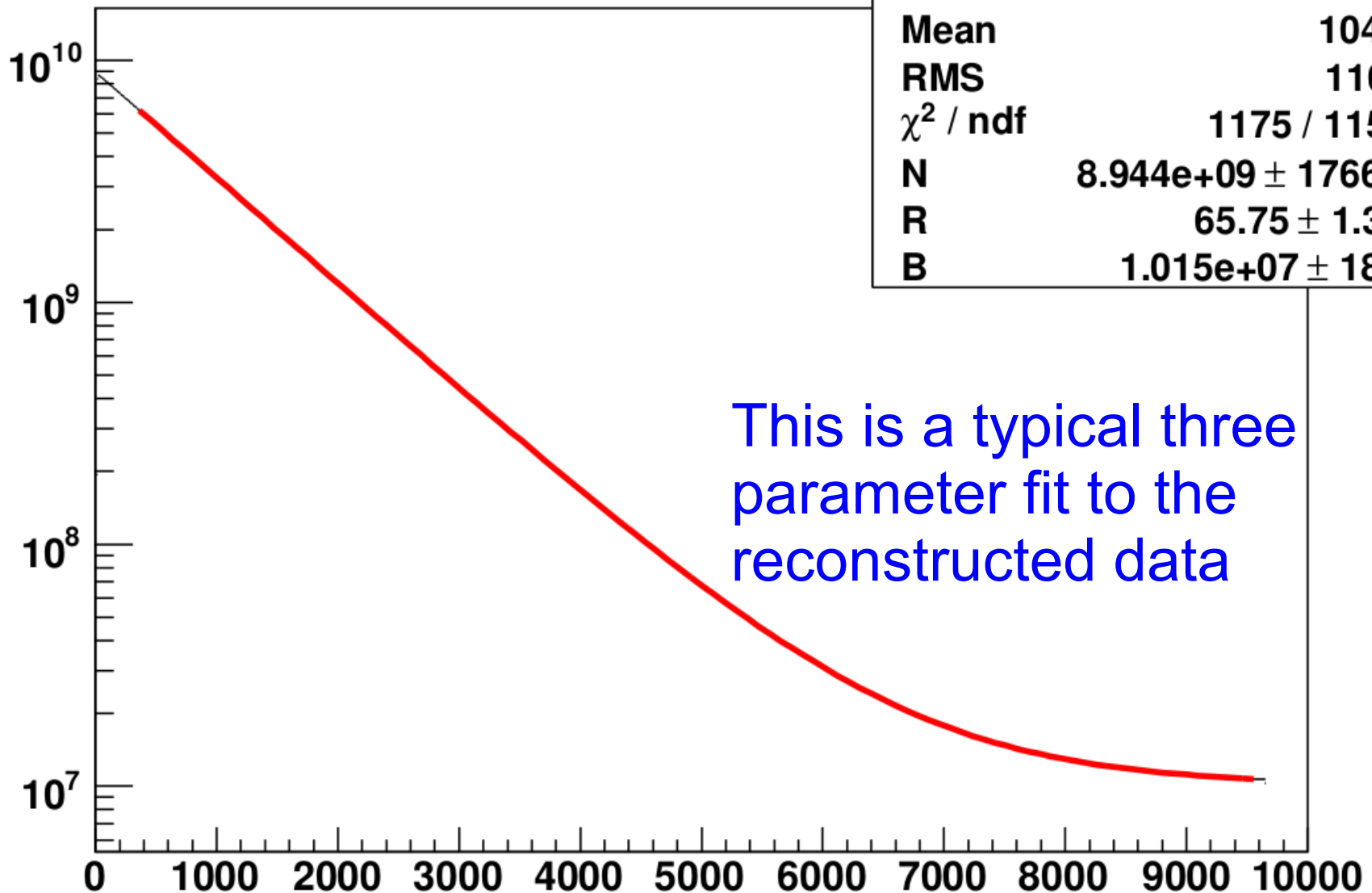
Nearly 10^{12} events on tape for each period...



... while new electronics and analysis techniques greatly reduced systematics over 2004.

Lifetime fits

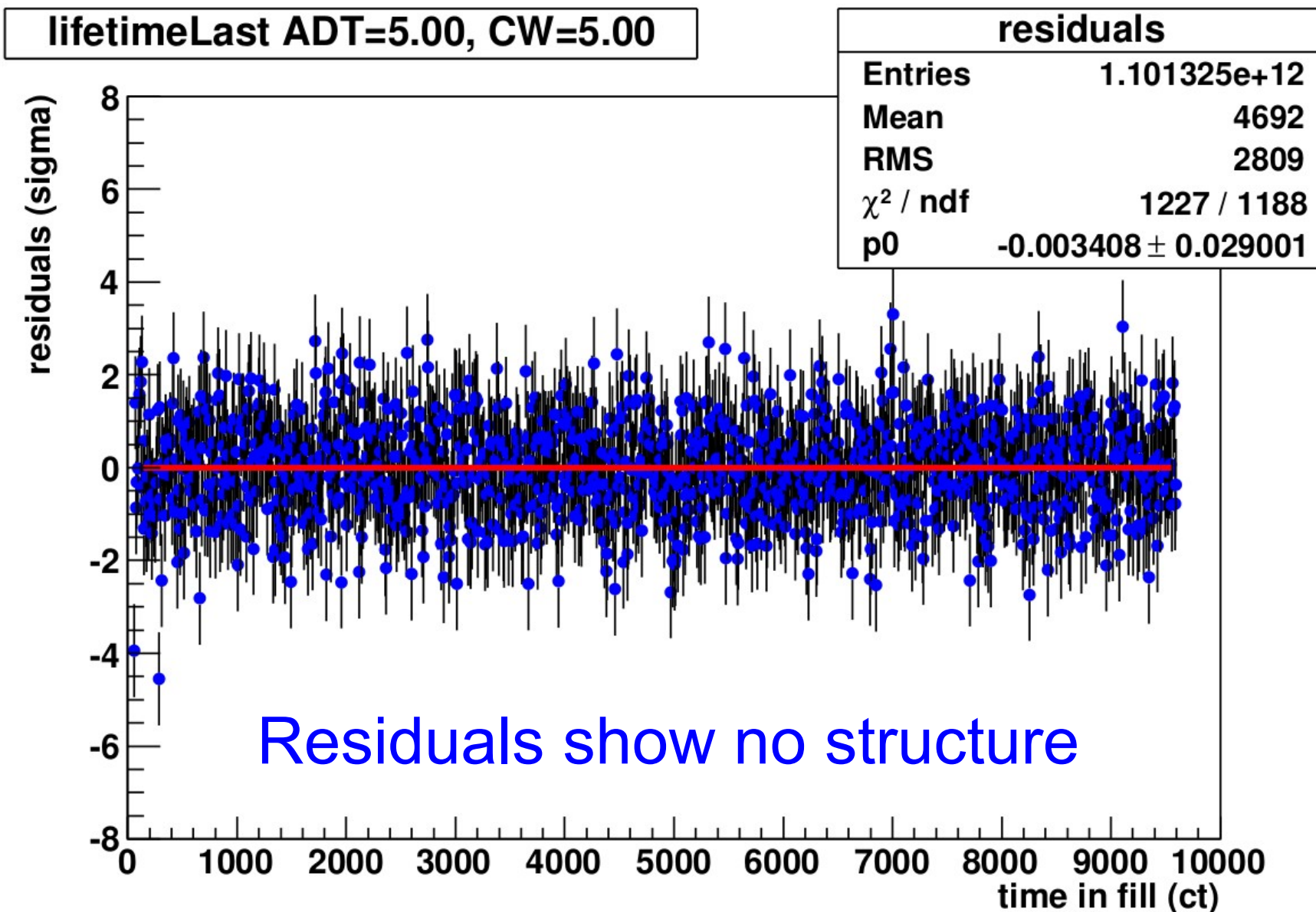
lifetimeLast ADT=8.00, CW=8.00



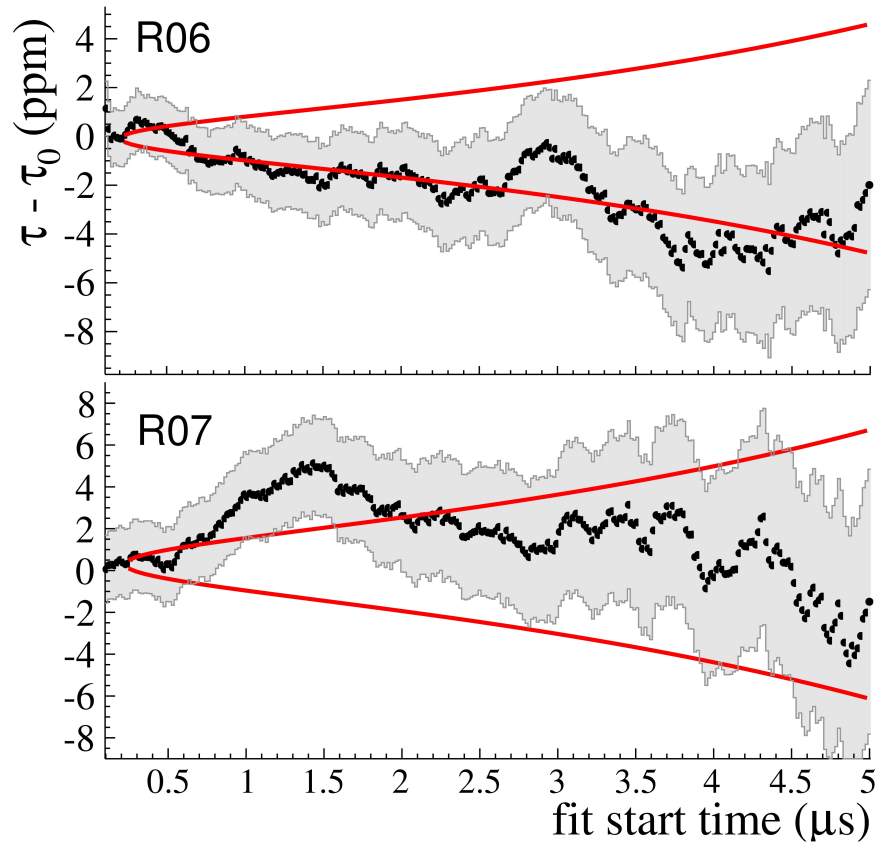
lifetimeLast4_px

| | |
|-----------------------|------------------------------|
| Entries | 1.101349e+12 |
| Mean | 1049 |
| RMS | 1101 |
| χ^2 / ndf | 1175 / 1158 |
| N | $8.944\text{e}+09 \pm 17662$ |
| R | 65.75 ± 1.30 |
| B | $1.015\text{e}+07 \pm 186$ |

Lifetime fits

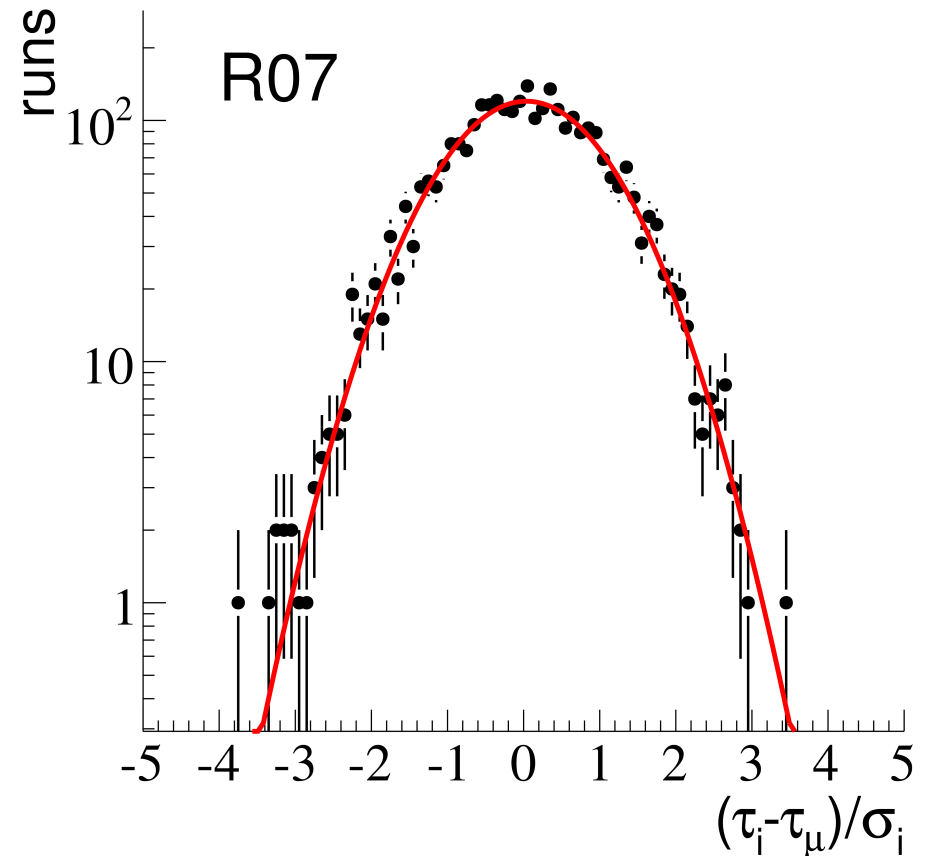


We studied a significant number of data subsets for lifetime consistency across various conditions



Fit start time scans show no evidence of missing long time scale components

The lifetimes measured by individual detector pairs appear statistically consistent



We believe our systematics are well understood for both run periods

| Uncertainty | R06 (ppm) | R07 (ppm) |
|---------------------------------|--------------|--------------|
| Kicker stability | 0.20 | 0.07 |
| μ SR distortions | 0.10 | 0.20 |
| Pulse pileup | 0.20 | |
| Gain variations | 0.25 | |
| Upstream stops | 0.10 | |
| Timing pick-off stability | 0.12 | |
| Master clock calibration | 0.03 | |
| Combined systematic uncertainty | 0.42 | 0.42 |
| Statistical uncertainty | 1.14 | 1.68 |

The lifetime results for our three run periods are entirely consistent

$$\tau_{\mu}^{\text{R06}} = 2\,196\,979.9 \pm 2.5(\text{stat}) \pm 0.9(\text{syst}) \text{ ps}$$

$$\tau_{\mu}^{\text{R07}} = 2\,196\,981.2 \pm 3.7(\text{stat}) \pm 0.9(\text{syst}) \text{ ps}$$

After properly accounting for the correlated systematics, the final combined MuLan result is

$$\tau_{\mu}^{\text{MuLan}} = 2\,196\,980.3 \pm 2.1(\text{stat}) \pm 0.7(\text{syst}) \text{ ps}$$

The lifetime results for our three run periods are entirely consistent

$$\tau_{\mu}^{\text{R06}} = 2\,196\,979.9 \pm 2.5(\text{stat}) \pm 0.9(\text{syst}) \text{ ps}$$

$$\tau_{\mu}^{\text{R07}} = 2\,196\,981.2 \pm 3.7(\text{stat}) \pm 0.9(\text{syst}) \text{ ps}$$

After properly accounting for the correlated systematics, the final combined MuLan result is

$$\tau_{\mu}^{\text{MuLan}} = 2\,196\,980.3 \pm 2.1(\text{stat}) \pm 0.7(\text{syst}) \text{ ps}$$

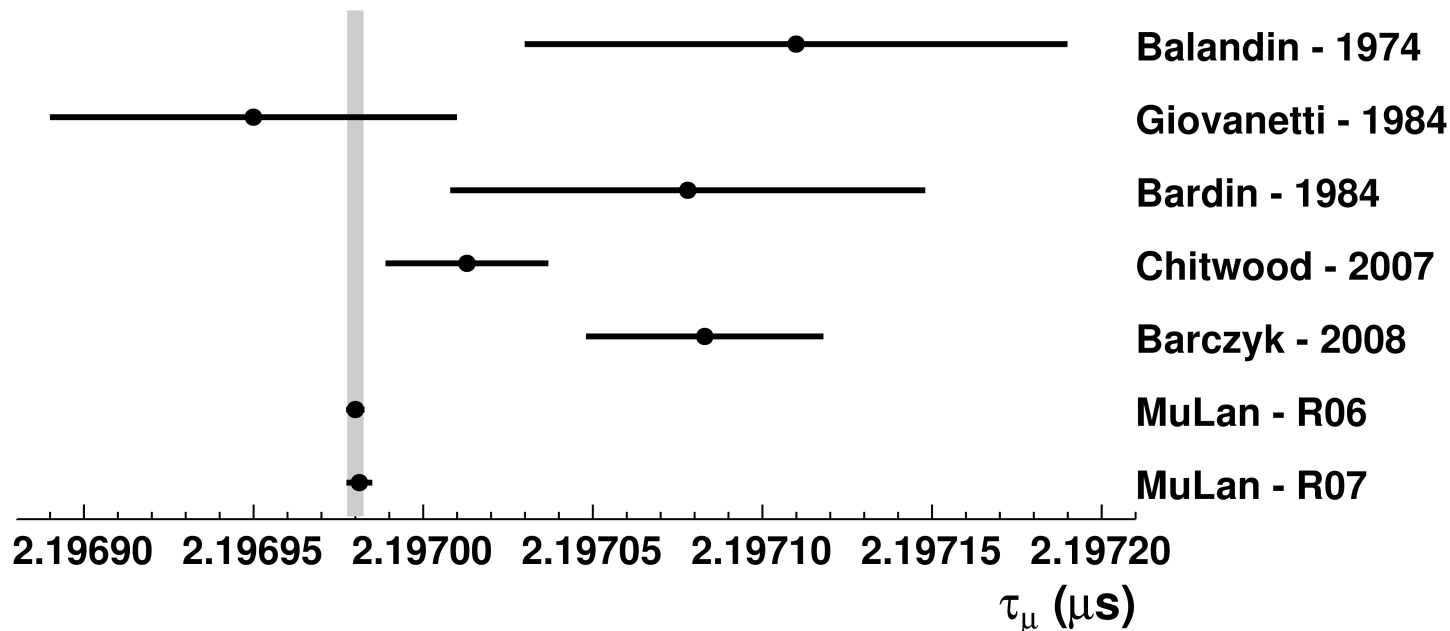

$$\pm 2.2 \text{ ps} \longrightarrow 1.0 \text{ ppm}$$

Our combined result dominates the world average

| Measured lifetime (μs) | Reference | Publication year |
|-------------------------------------|------------|------------------|
| $2.196\,9803 \pm 0.000\,0022$ | R06+R07 | |
| $2.196\,9799 \pm 0.000\,0027$ | R06 | |
| $2.196\,9812 \pm 0.000\,0038$ | R07 | |
| $2.197\,083 \pm 0.000\,035$ | Barczyk | 2008 |
| $2.197\,013 \pm 0.000\,024$ | Chitwood | 2007 |
| $2.197\,078 \pm 0.000\,073$ | Bardin | 1984 |
| $2.196\,95 \pm 0.000\,06$ | Giovanetti | 1984 |
| $2.197\,11 \pm 0.000\,08$ | Balandin | 1974 |
| $2.197\,3 \pm 0.000\,3$ | Duclos | 1973 |

$$\tau_{\mu}^{\text{PDG}} = 2\,196\,981.1 \pm 2.2 \text{ ps}$$

There is some tension with the Barczyk result (FAST) that drives the increased error bar in the PDG average.



Our motivation, of course, is extracting the Fermi Constant

Assuming a pure V-A structure of the weak interactions, we can extract Fermi's constant by inverting:

$$\frac{1}{\tau_\mu} = \frac{G_F^2 m_\mu^5}{192\pi^3} \left(1 + \Delta q^{(0)} + \Delta q^{(1)} + \Delta q^{(2)} \right)$$

Phase space $\rightarrow \Delta q^{(0)}$

First order corrections $\rightarrow \Delta q^{(1)}$

Second order corrections $\rightarrow \Delta q^{(2)}$

$$G_F^{\text{MuLan}} = 1.166\,378\,7(6) \times 10^{-5} \text{ GeV}^{-2}$$

$\rightarrow 0.5 \text{ ppm}$

MuLan was systematics limited ... could we do better at a future facility?

| Uncertainty | R06 (ppm) | R07 (ppm) |
|---------------------------------|--------------|--------------|
| Kicker stability | 0.20 | 0.07 |
| μ SR distortions | 0.10 | 0.20 |
| Pulse pileup | 0.20 | |
| Gain variations | 0.25 | |
| Upstream stops | 0.10 | |
| Timing pick-off stability | 0.12 | |
| Master clock calibration | 0.03 | |
| Combined systematic uncertainty | 0.42 | 0.42 |
| Statistical uncertainty | 1.14 | 1.68 |

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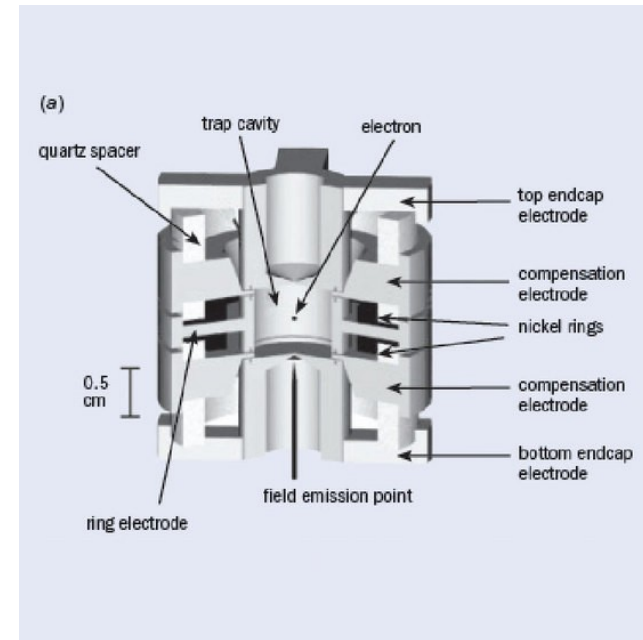
My Verdict: Probably ...

Precision electroweak parameters: an update

Fine Structure Constant

$$\frac{\delta\alpha_{\text{em}}}{\alpha_{\text{em}}} \approx 0.32 \text{ ppb}$$

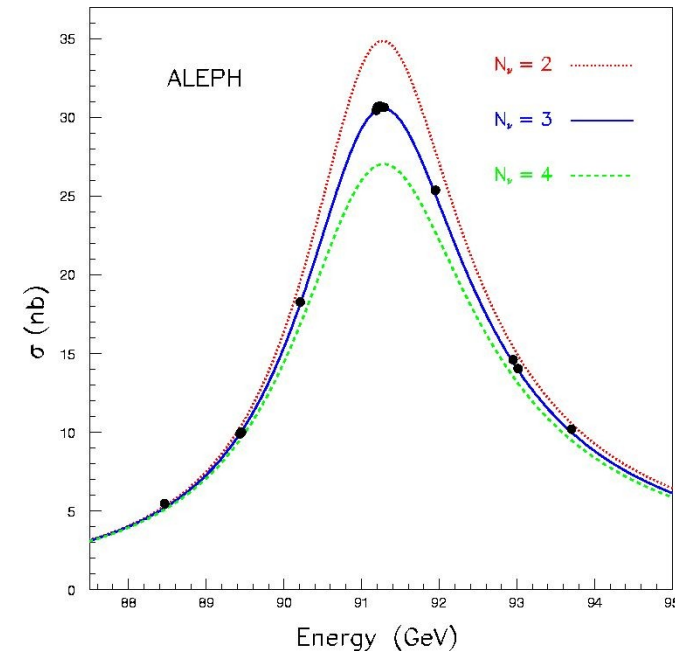
Gabrielse *et al*
2008



Mass of the neutral weak boson

$$\frac{\delta M_{Z^0}}{M_{Z^0}} \approx 23 \text{ ppm}$$

LEP EWWG
2005

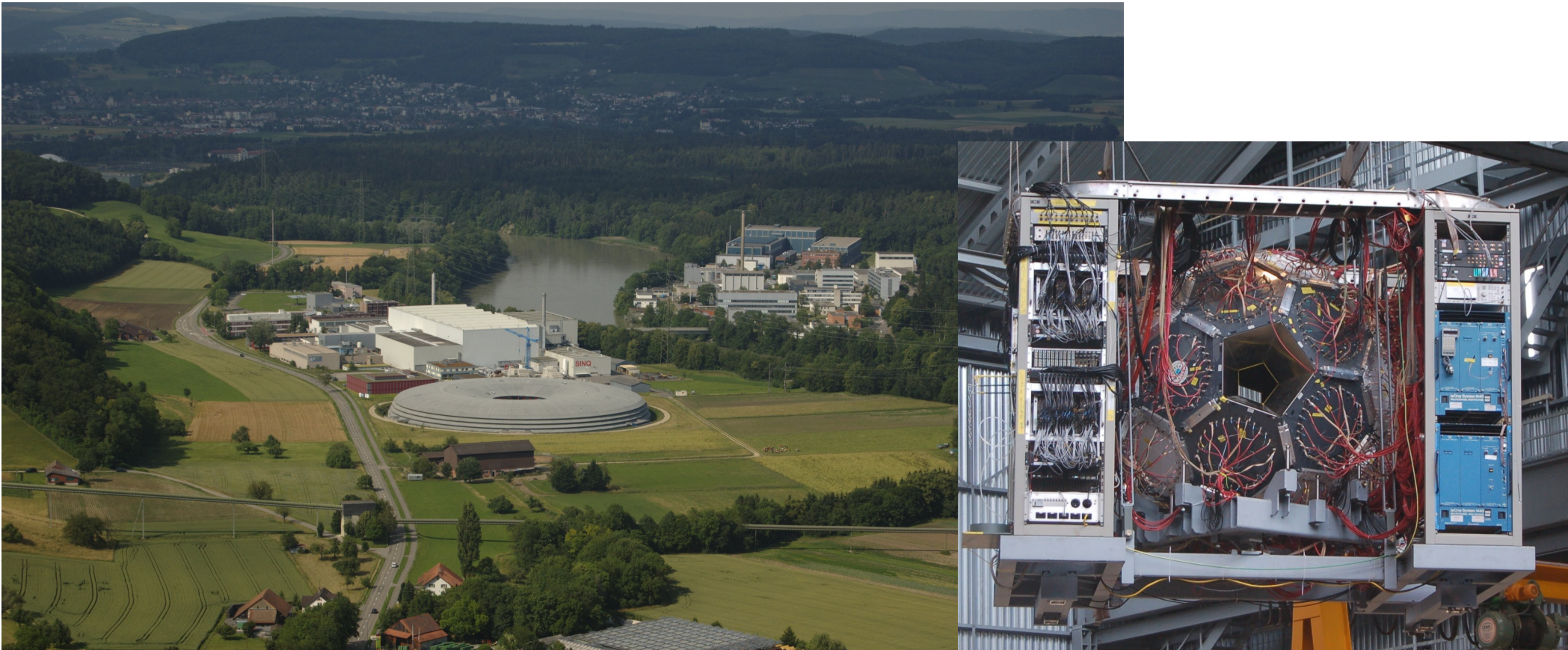


Fermi Constant

$$\frac{\delta G_F}{G_F} \approx 0.5 \text{ ppm}$$

Tishchenko, et al
2012

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Residual polarization of the stopped muons plays havoc with lifetime fits

Weak decay
violates chirality

$$\frac{d\Gamma^\pm}{d(\cos \theta)} = \frac{1}{\tau_\mu} \left(1 \pm \frac{1}{3} A \cos \theta \right)$$

$$N_D(t) = N_0 e^{-t/\tau} \left(1 + \frac{1}{3} A \left[\vec{S}_\perp(t) \cdot \hat{e}_D e^{-t/T_\perp} + \vec{S}_\parallel(0) \cdot \hat{e}_D e^{-t/T_\parallel} \right] \right)$$

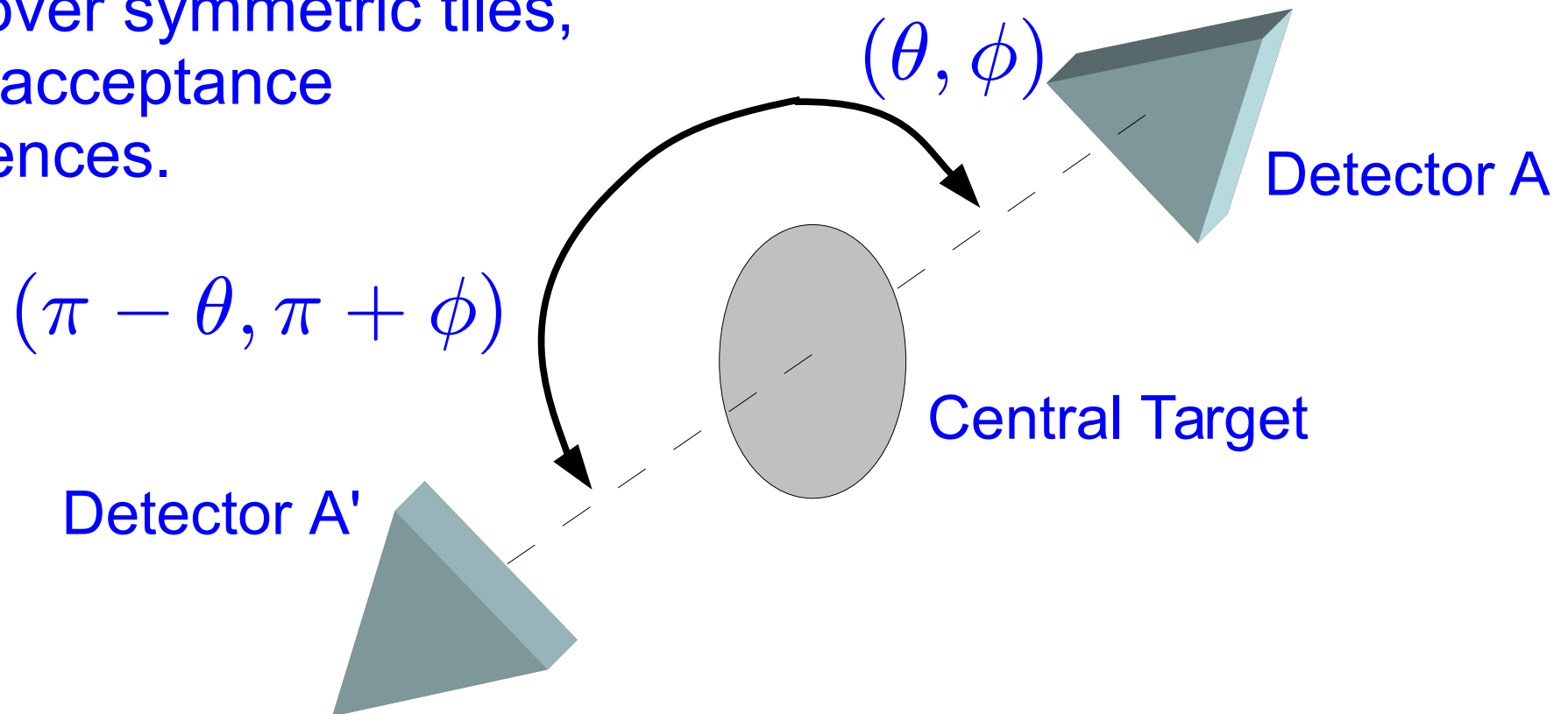
Spins precess in
magnetic fields

$$H = -\vec{\mu} \cdot \vec{B}$$

Matter interactions decrease
polarization fraction over
time

We start with nearly 100% polarized beam ...
how do we control polarization issues?

Point symmetry of the
detector largely cancels
polarization asymmetries in
sum over symmetric tiles,
up to acceptance
differences.



We also modulate the remnant polarization by choice of target environment and muonium formation fraction

A polarization destroying ferromagnetic target, AK3, with high internal field (2004,2006)

Polarization preserving target, crystalline quartz, with an applied external field (2007)

We also performed special runs with polarization maximizing targets like copper and aluminum, and target offsets to maximize asymmetries (2006, 2007)